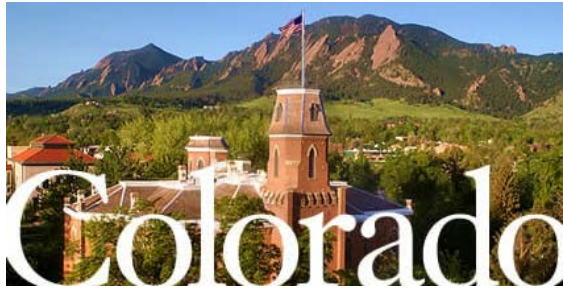


Experiments and Model Development for Polydisperse, Gas-fluidized Systems



Jia-Wei Chew
R. Brent Rice
Christine M. Hrenya

Vicente Garzó

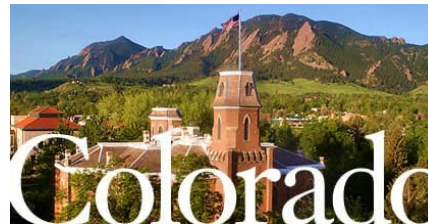
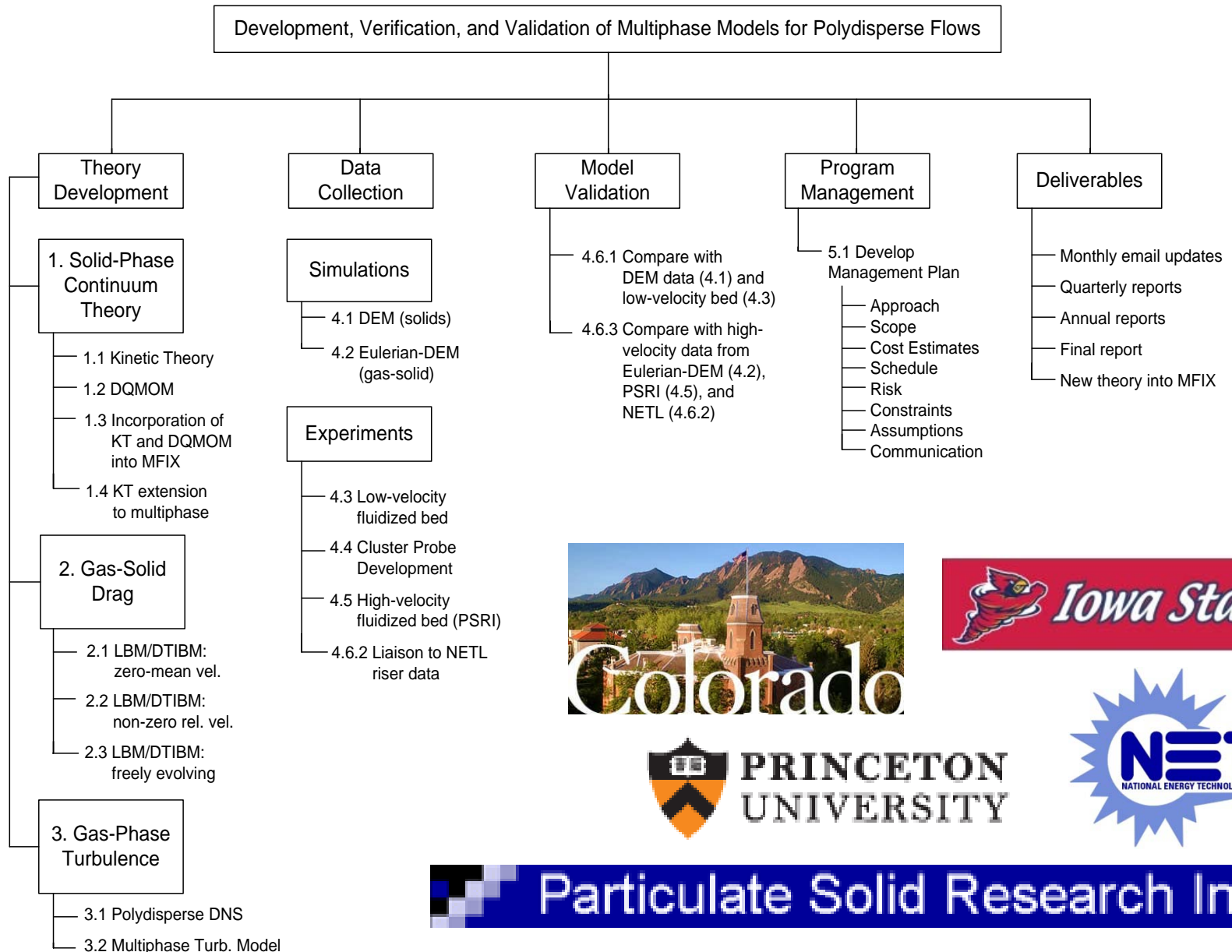


S. Tenneti
R. Garg
Shankar Subramaniam

NETL 2010 Workshop on Multiphase Flow Science
4 May 2010
Pittsburgh, PA



Project Scope: Work Breakdown Structure



Univ. Colorado Tasks for Year 3

Validation data

- 1) **Bubbling-bed experiments (Colorado)**
- 2) **DEM simulations of simple shear (Colorado)**
- 3) **Riser experiments (Colorado & PSRI)**

Ray Cocco (PSRI) – Thursday

Theory

- 4) **Application of polydisperse kinetic theory to clustering (Princeton & NETL & Colorado)**

Bill Holloway (Princeton) - Tuesday

- 5) **Kinetic Theory Extension to Gas-Solid Flows (Colorado & Iowa State)**

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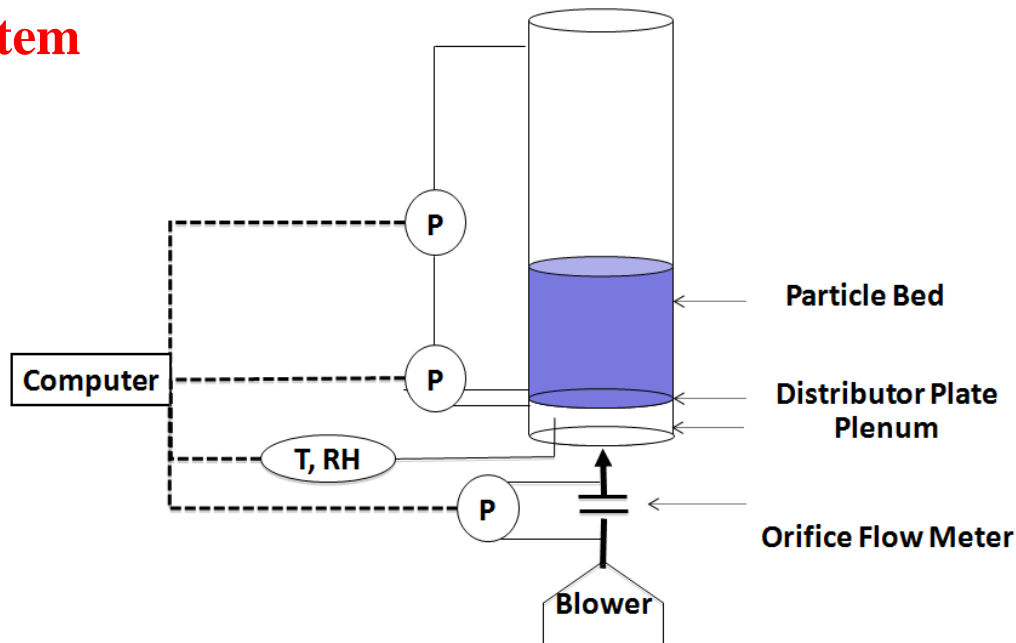
Overview: Bubbling Bed Experiments

Objective

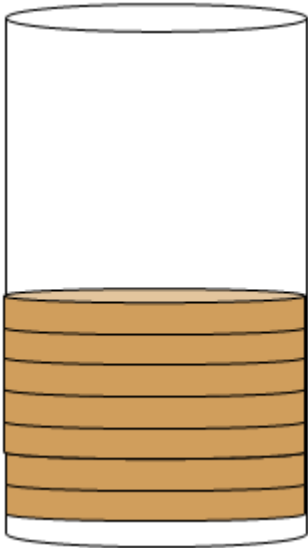
To characterize *segregation* and *bubbling* behavior of bubbling beds with *continuous size distributions*

- Do continuous PSD's behave like binary mixtures?
- Is there a direct link between the segregation and bubbling levels?

System

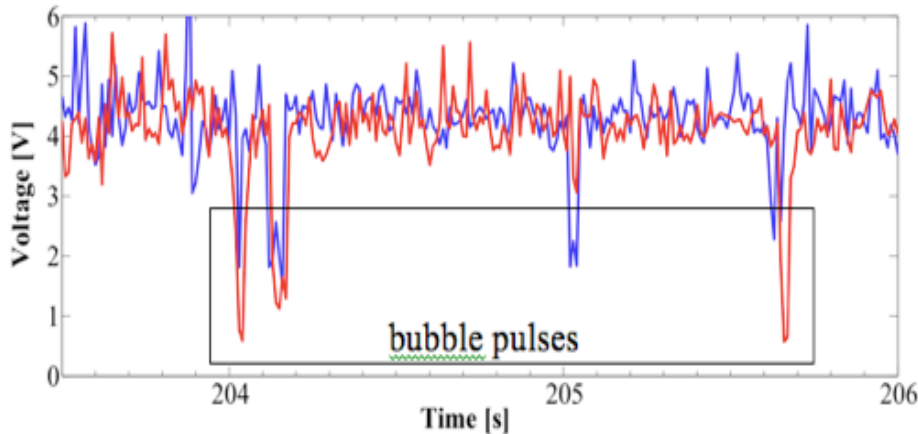
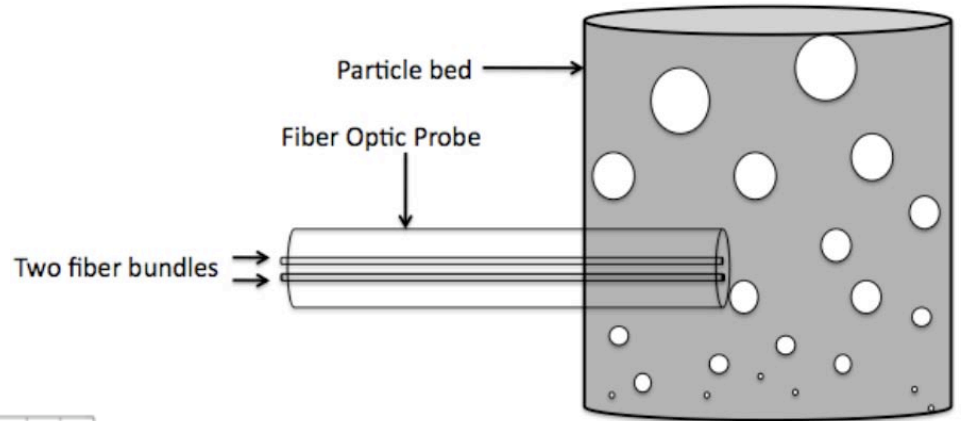


Measurements: Segregation and Bubbling



Segregation: sieving of thin vertical sections

1. High velocity ($3 U_{cf}$) to mix bed
2. Low velocity ($1.2 U_{cf}$) for segregation
3. Shut down, vacuum sections, sieve

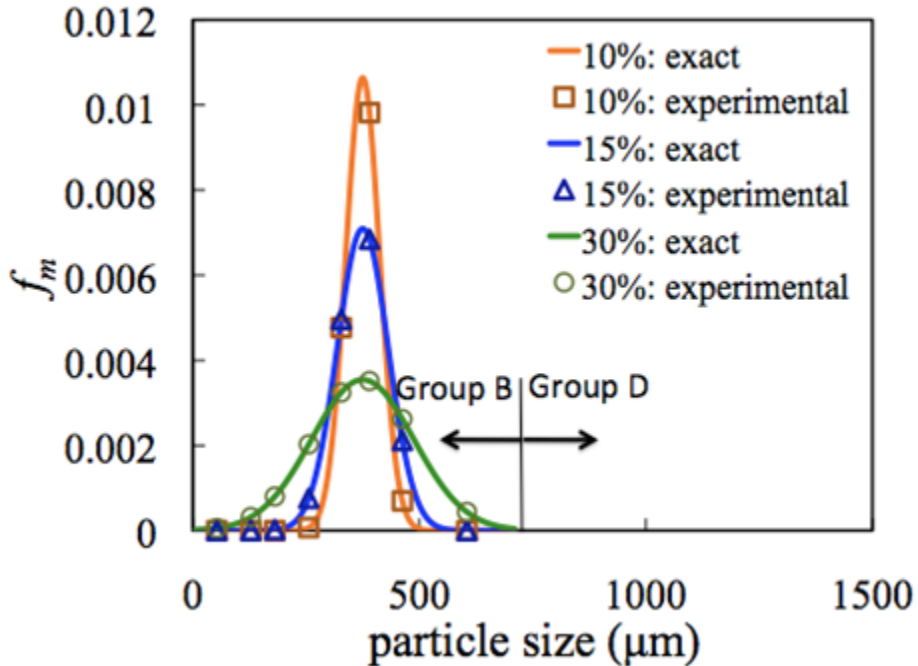


Bubbling: fiber optic probe

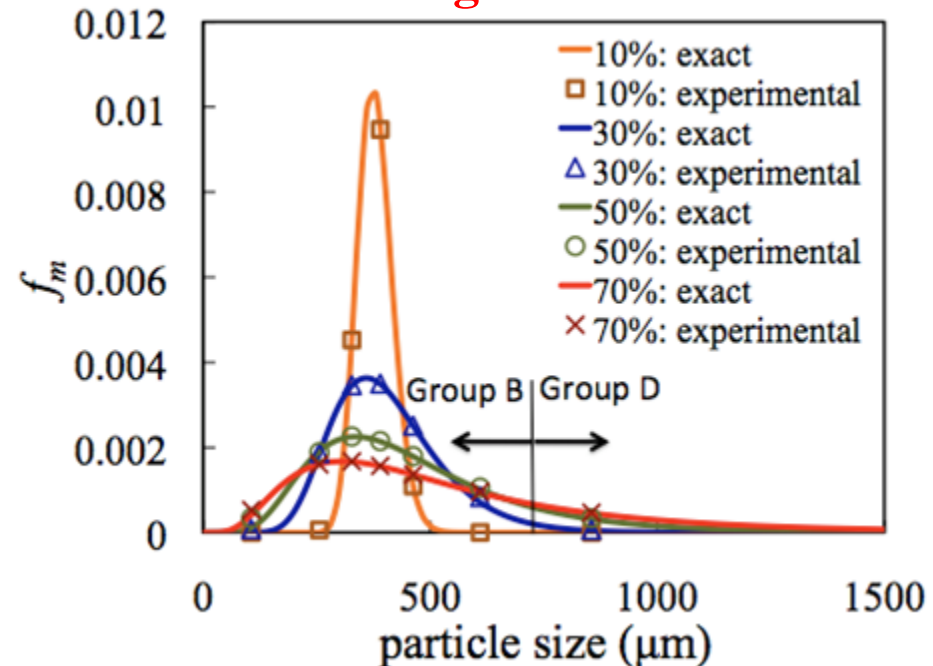
1. 7 axial x 9 radial positions
2. Bubble frequency, velocity & size

Continuous PSD's Investigated

Gaussian



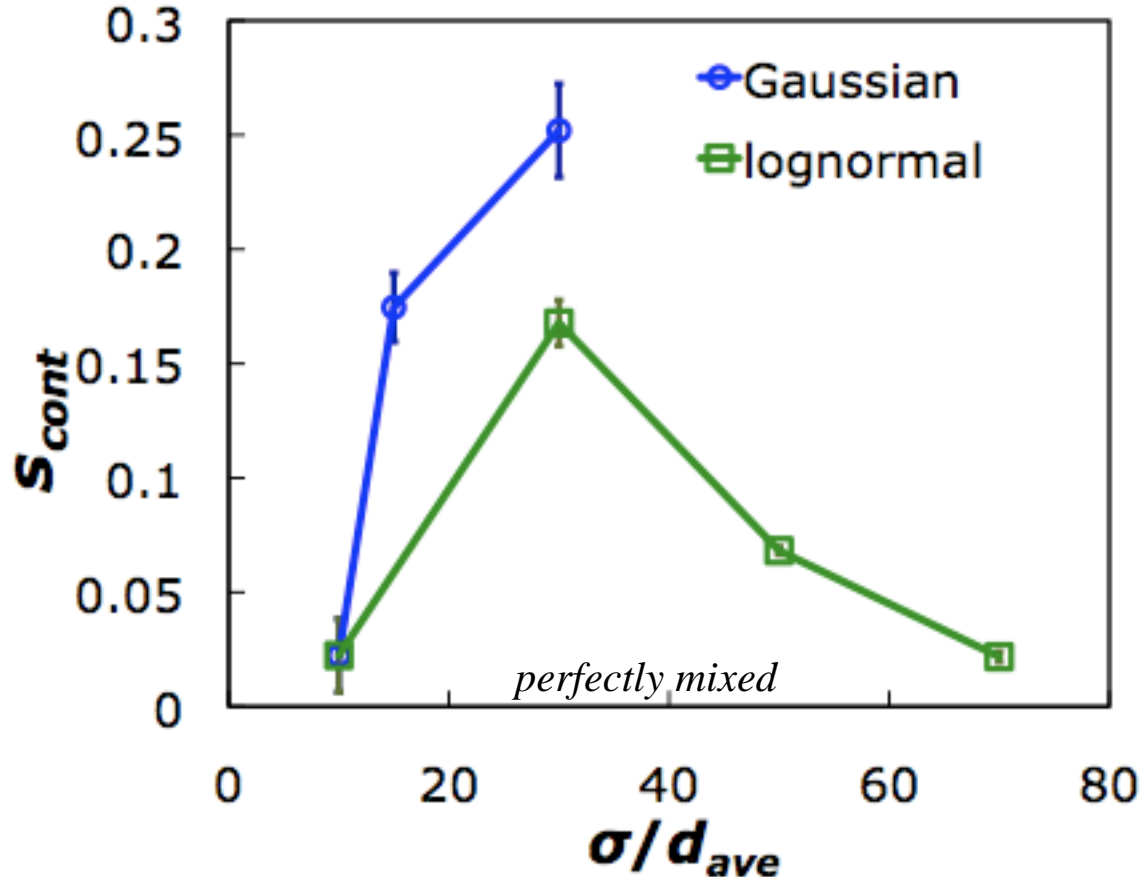
lognormal



Width of distribution: σ/d_{ave} (%)

Material: sand

Results: Segregation



$$S_{cont} = \frac{S - 1}{S_{max} - 1}$$

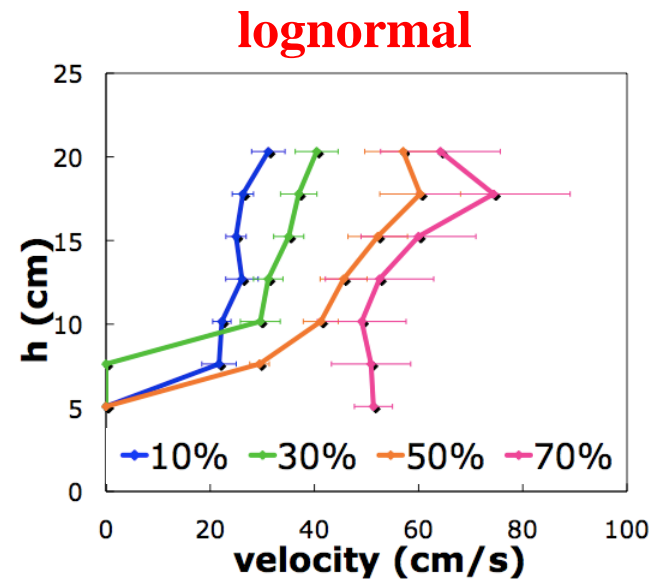
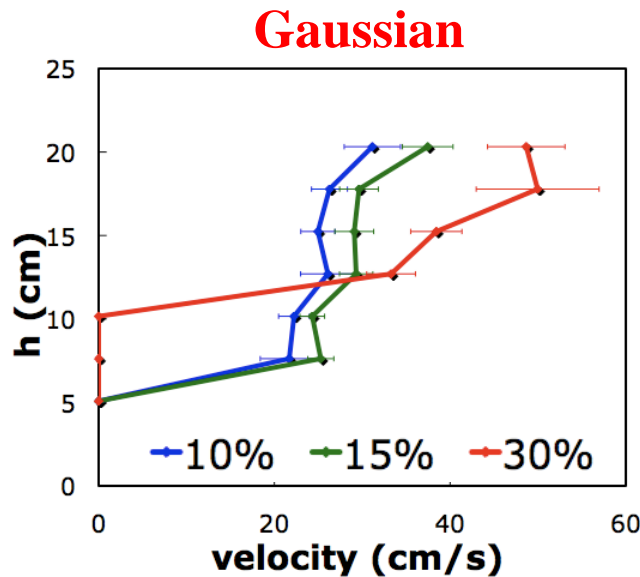
$$S = \frac{\langle h_{small} \rangle}{\langle h_{large} \rangle} \quad S_{max} = \frac{2x_{large} + x_{small}}{x_{large}}$$

$S_{cont} = 0 \rightarrow$ perfect mixing
 $S_{cont} = 1 \rightarrow$ perfect segregation

Gaussian: as PSD width increases \uparrow , segregation \uparrow

lognormal: *non-monotonic* variation of segregation with PSD width

Results: Bubbling

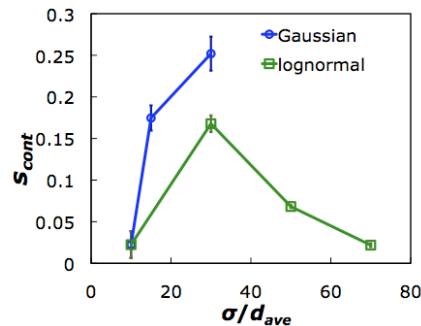
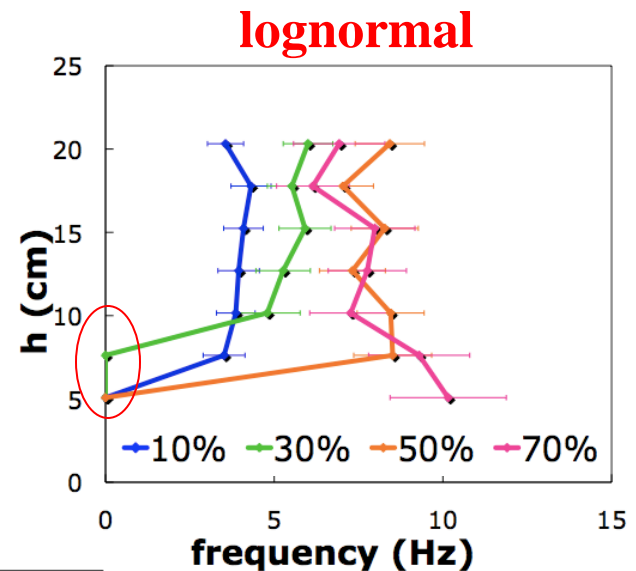
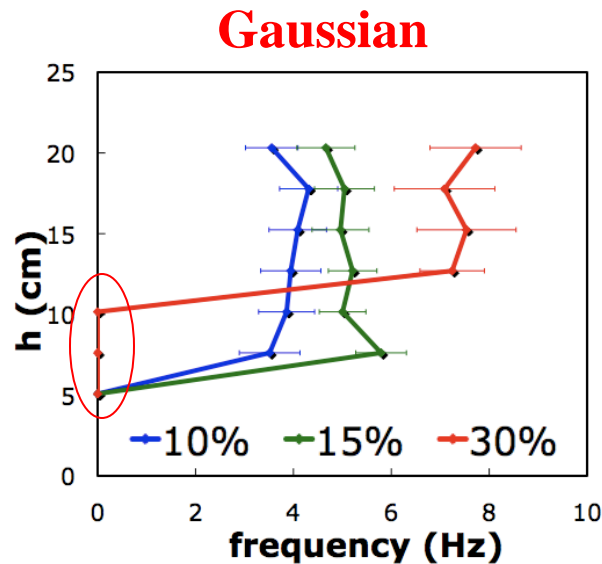


Gaussian: as PSD width increases \uparrow , all bubble parameters \uparrow (some not shown)
lognormal: same as Gaussian (*monotonic* variation)

Reconciliation: Segregation and Bubbling Measurements

Explanation

- presence of *bubble-less layer* in some systems
- size of bubble-less layer correlates with degree of segregation



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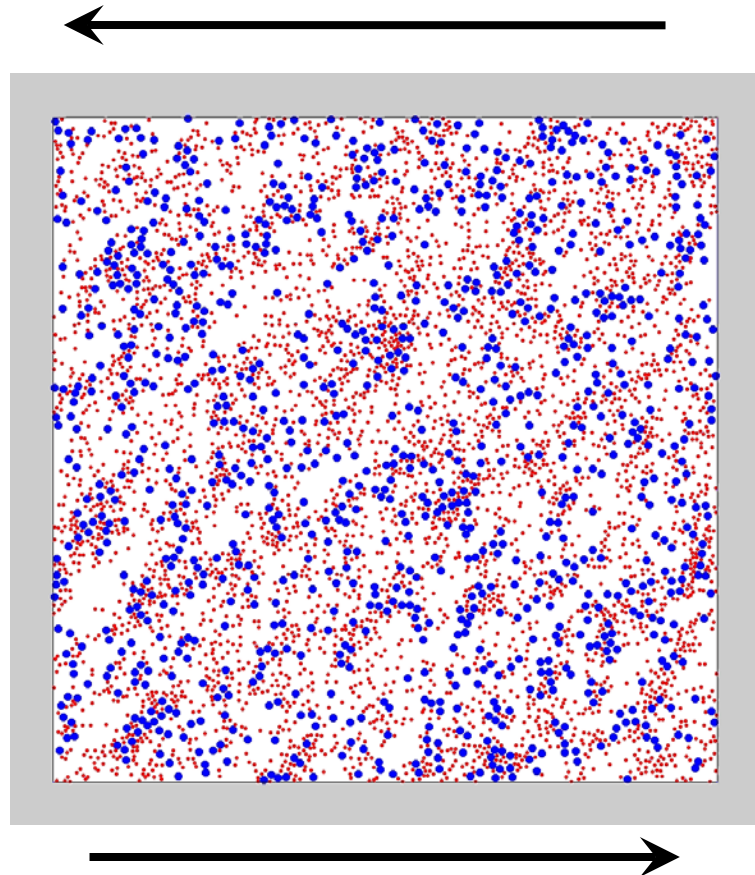
Objective

Assess the impact of *polydispersity* on *clustering* in *granular* flows

- How does polydispersity affect the prominence of clusters?
- Does species segregation occur in a *transient* cluster?

System: Simple Shear Flow

Approach: MD simulations

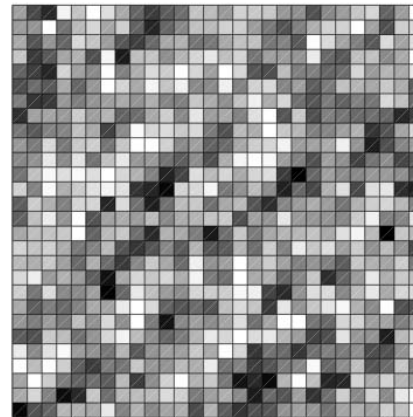
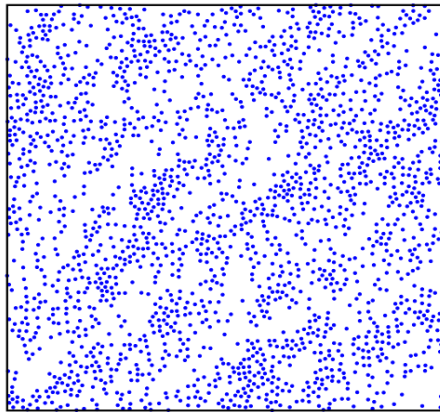


MD Simulations

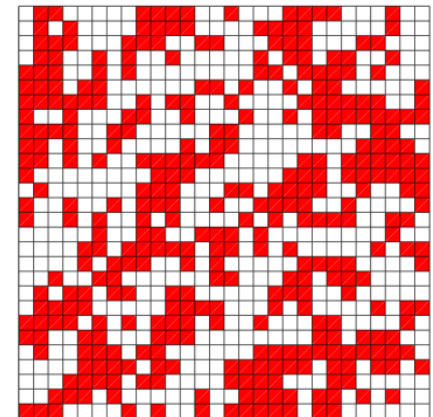
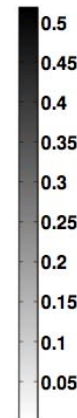
Simulation Description

- 2D, event-driven
- Inelastic, frictionless, hard disks
- Size distributions: *binary, Gaussian, lognormal*

Concentration mapping



$$v_{avg} = 0.2$$



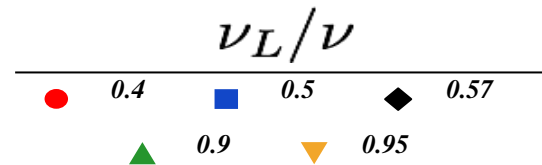
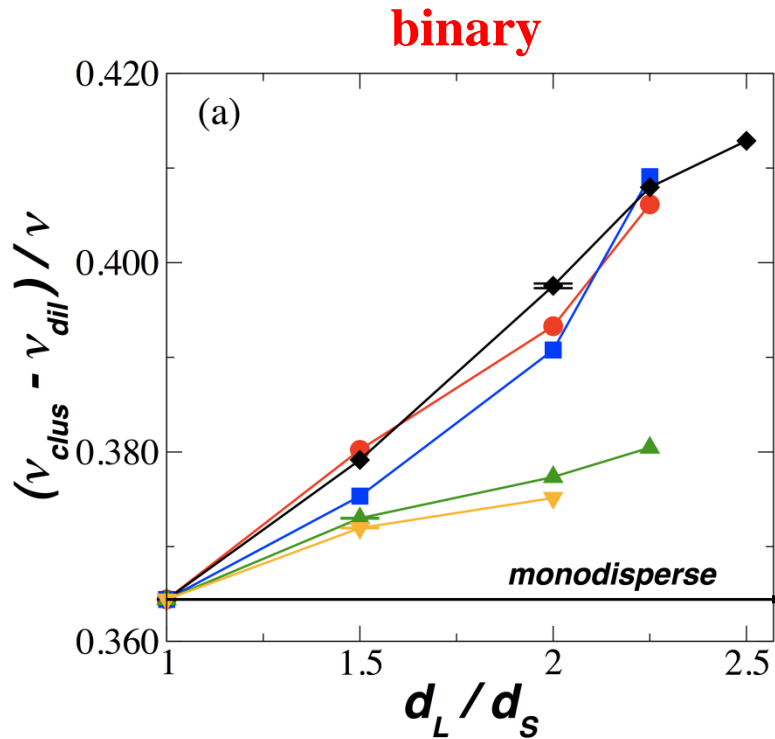
Cluster Region: $v > v_{avg}$

Dilute Region: $v < v_{avg}$

Resulting Quantities: v_{clus} v_{dil} T_{clus} T_{dil}

$v_{L,clus}$ $v_{L,dil}$ $T_{L,clus}$ $T_{L,dil}$

Results: Cluster Prominence



Increased Prominence
Greater Conc. Difference

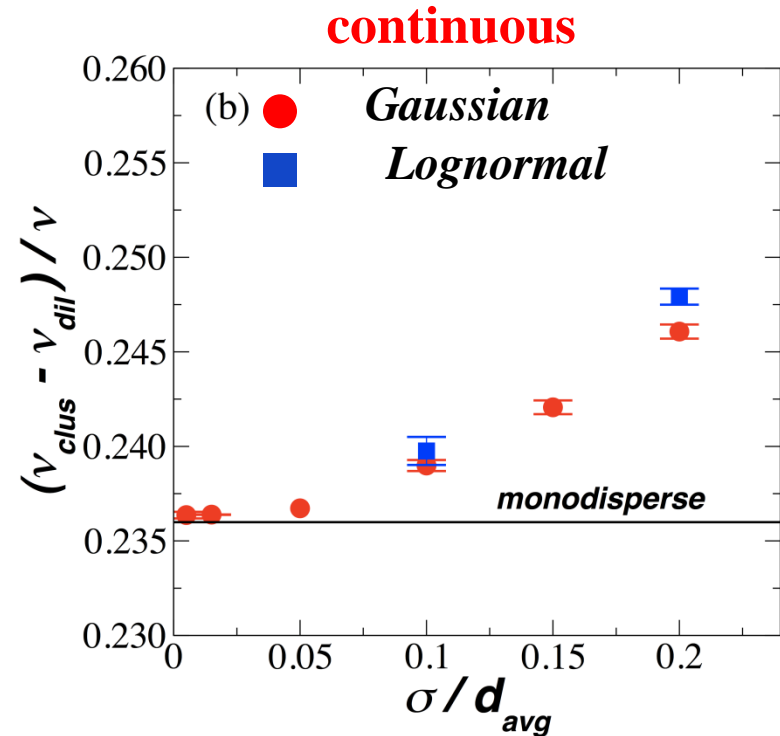
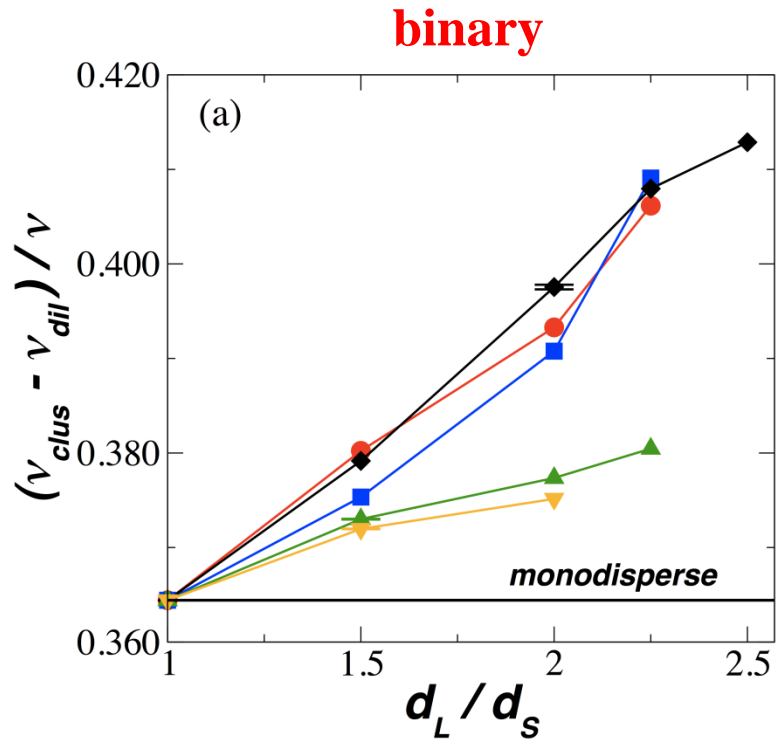
Lesser Conc. Difference
Decreased Prominence

Cluster prominence greater for systems with more than one species



Tendency increases with deviation from monodisperse limit

Results: Cluster Prominence

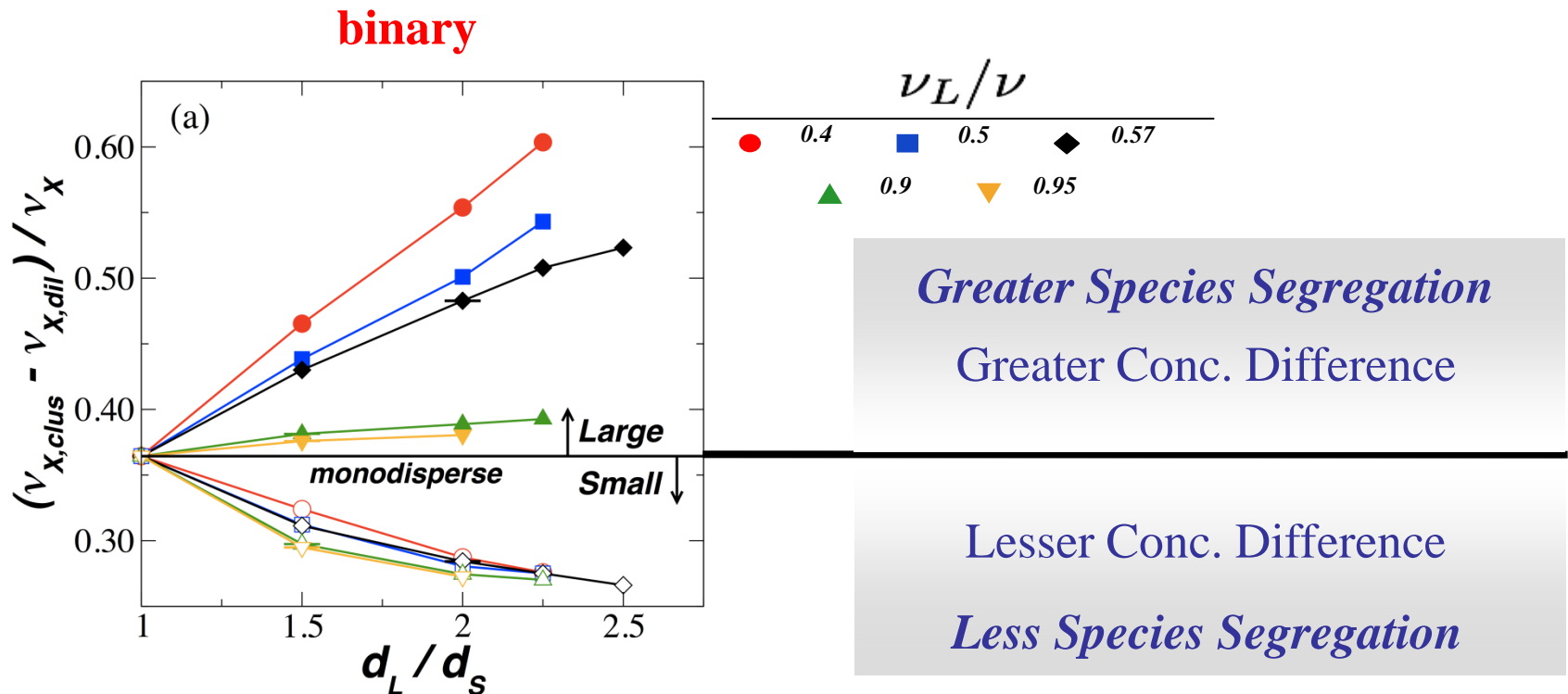


Cluster prominence greater for systems with more than one species



Tendency increases with deviation from monodisperse limit

Results: Species Segregation

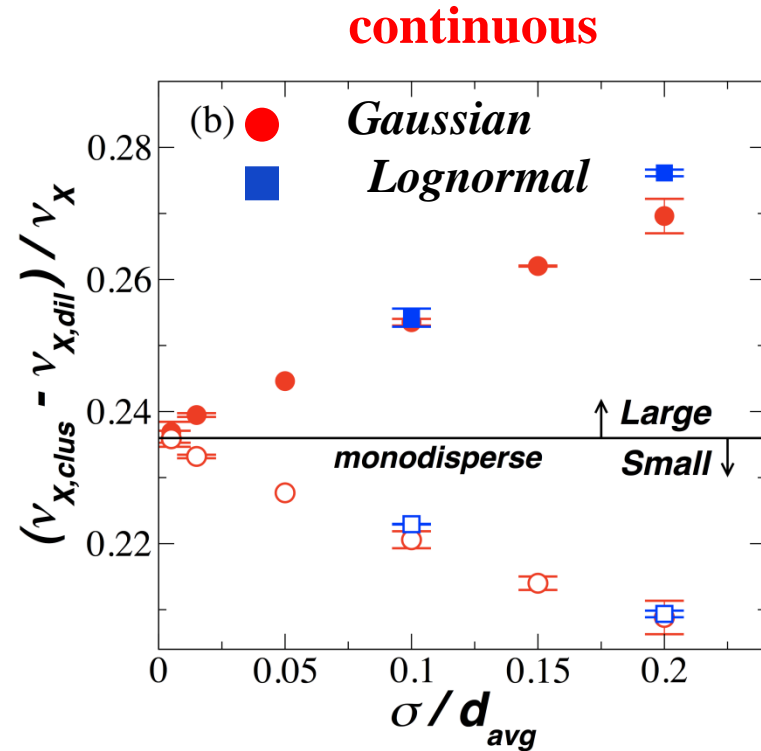
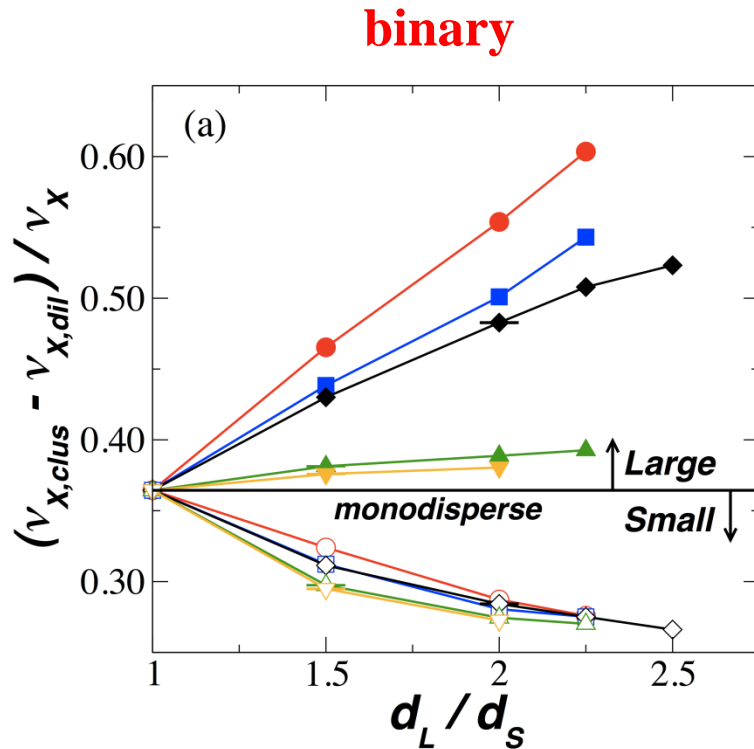


Large Particles segregate preferentially toward the clustered regions



Tendency increases with increasing size disparity

Results: Species Segregation



Large Particles segregate preferentially toward the clustered regions



Tendency increases with increasing size disparity

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Modeling of Gas-solid Flows

Continuum (“Two-fluid”) Description

- Gas phase: Navier-Stokes + turbulence + drag force
- Solids phase: Kinetic-theory-based models + fluid-phase interactions

↑
*modifications to
kinetic-theory closures*

↑
new terms

Current Objective: *Incorporation of gas-phase (drag) effects into kinetic-theory-based models for solid phase*

Physical Picture

Recall fluid-solid interaction force (drag force)

$$F_{fluid} = F_n + F_t = \int_0^{2\pi} \int_0^{\pi} \left(-p|_{r=R} \cos \theta \right) R^2 \sin \theta d\theta d\phi + \int_0^{2\pi} \int_0^{\pi} \left(\tau_{r\theta}|_{r=R} \sin \theta \right) R^2 \sin \theta d\theta d\phi$$

$\left(-p|_{r=R} \cos \theta \right)$
 $\left(\tau_{r\theta}|_{r=R} \sin \theta \right)$

→ = f (velocity & pressure at surface)

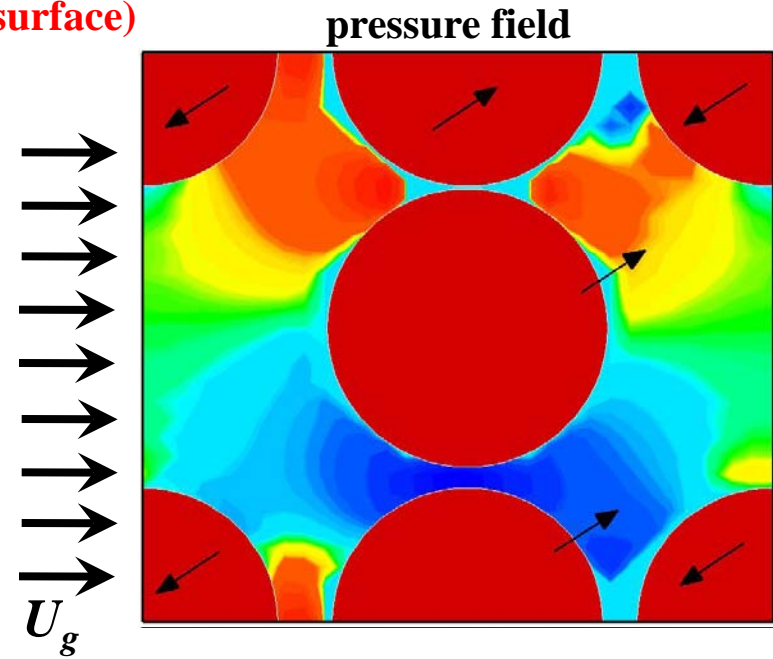
Mean fluid force on *single* particle

$$F_{fluid} = f(U_g - U)$$

Velocity & pressure fields (& thus fluid force) change with:

- Fluctuations in particle velocity
- Fluctuations in gas velocity

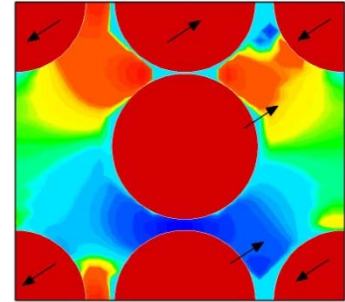
→ Mean vs. Instantaneous Fluid Force



Incorporation of Instantaneous Fluid Force

Alternative 1: DEM (solids) /DNS (fluid) – resolve flow field around particles

- + fluid force is “output”
- too computationally expensive
(no-slip BC at each moving particle surface)



Alternative 2: Two-fluid model – “averaged” flow field over several particles

- + computationally feasible (single equation of motion for each phase)
- fluid effects are “input” – model is needed to subsume *instantaneous* effects

→ Q1: Impact on governing equations (additional terms)?

→ Q2: Impact on constitutive relations for solid phase (\mathbf{P} , \mathbf{q} , ζ)?

Current Approach

→ (i) use DEM/DNS simulations to develop model for instantaneous force

→ (ii) incorporate this force model into starting kinetic equation & derive hydrodynamic description ←

More details...

Basic Idea: Incorporation of fluid force into Enskog kinetic equation

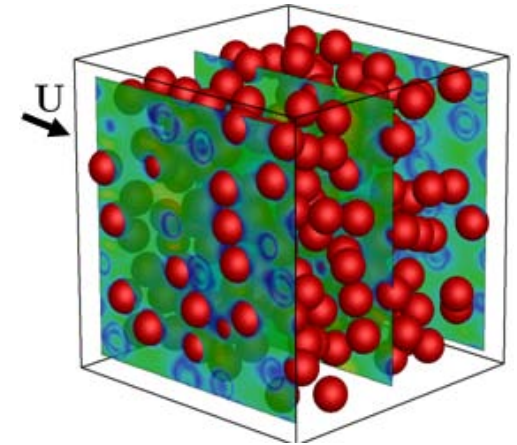
$$\frac{\partial}{\partial t} f + v_i \frac{\partial f}{\partial x_i} + \frac{\partial}{\partial v_i} \left(\frac{F_{fluid,i}}{m} \right) + g_i \frac{\partial}{\partial v_i} f = J$$

transient convective $\frac{\partial}{\partial v_i} \left(\frac{F_{fluid,i}}{m} \right)$ *gravity* *collisions*

*instantaneous fluid force
on single particle*

**DEM/DNS technique for closure: IBM (Immersed Boundary Method)
based model of $F_{fluid,i}$ as function of:**

- Hydrodynamic variables: ϕ , U_i , T , U_{gi}
- Physical parameters: m , d , α , μ_g , ρ_g



IBM-based model for acceleration

$$\frac{F_{fluid,i}}{m} = A_i = -\frac{\beta_{IBM}}{m} (U_i - U_{gi}) - \frac{1}{m} \gamma_{ij} V_j - B_{ij} dW_j$$

coefficients extracted from IBM simulations

instantaneous particle acceleration
 mean particle velocity
 mean gas velocity
 fluctuating particle velocity
 Wiener process increment (stochastic model for fluctuating fluid velocity)

Use IBM simulations to find β^* , γ_{ij}^* , and B_{ij}^* as functions of

- ϕ solids volume fraction
- ρ_s/ρ_f density ratio
- $Re_m = \frac{(1-\phi)\rho_g d |\mathbf{U} - \mathbf{U}_{\infty}|}{\mu_g}$ particle Re based on mean flow
- $Re_T = \frac{\rho_g d}{\mu_g} \sqrt{\frac{T}{m}}$ particle Re based on particle velocity fluctuations

Resulting Hydrodynamic Description

Balance Equations (Solid-Phase Momentum & Granular Energy)

$$D_t \mathbf{U} + \frac{1}{mn} \nabla \mathbf{P} = \underbrace{-\frac{\beta_{IBM}}{m} (\mathbf{U} - \mathbf{U}_g)}_{\text{mean drag}} + \mathbf{g}$$

$$D_t T + \frac{2}{3n} (\nabla \cdot \mathbf{q} + P_{ij} \nabla_j U_i) = -\zeta T \underbrace{-\frac{2}{3\rho} \gamma_{ij} P_{ij}^k}_{\text{sink due to viscous drag}} + \underbrace{\frac{\rho}{3n} B_{ij} B_{ij}}_{\text{source due to fluid-particle fluctuations}}$$

Explicit Constitutive relations obtained for ζ , \mathbf{P} , and \mathbf{q} :

- Cooling rate $\zeta^{(0)} = \zeta^{(0)} (B_{ij})$
- Cooling rate TC $\zeta_U = \zeta_U (\gamma_{ij}, B_{ij})$
- Shear viscosity $\eta = \eta (\gamma_{ij}, B_{ij})$
- Bulk viscosity $\lambda = \lambda (B_{ij})$
- Conductivity $\kappa = \kappa (\gamma_{ij}, B_{ij})$
- Dufour coefficient $\mu = \mu (\gamma_{ij}, B_{ij})$

Base Case: Massive Particles ($St \gg 1$) and Stokes flow ($Re_m \ll 1$)

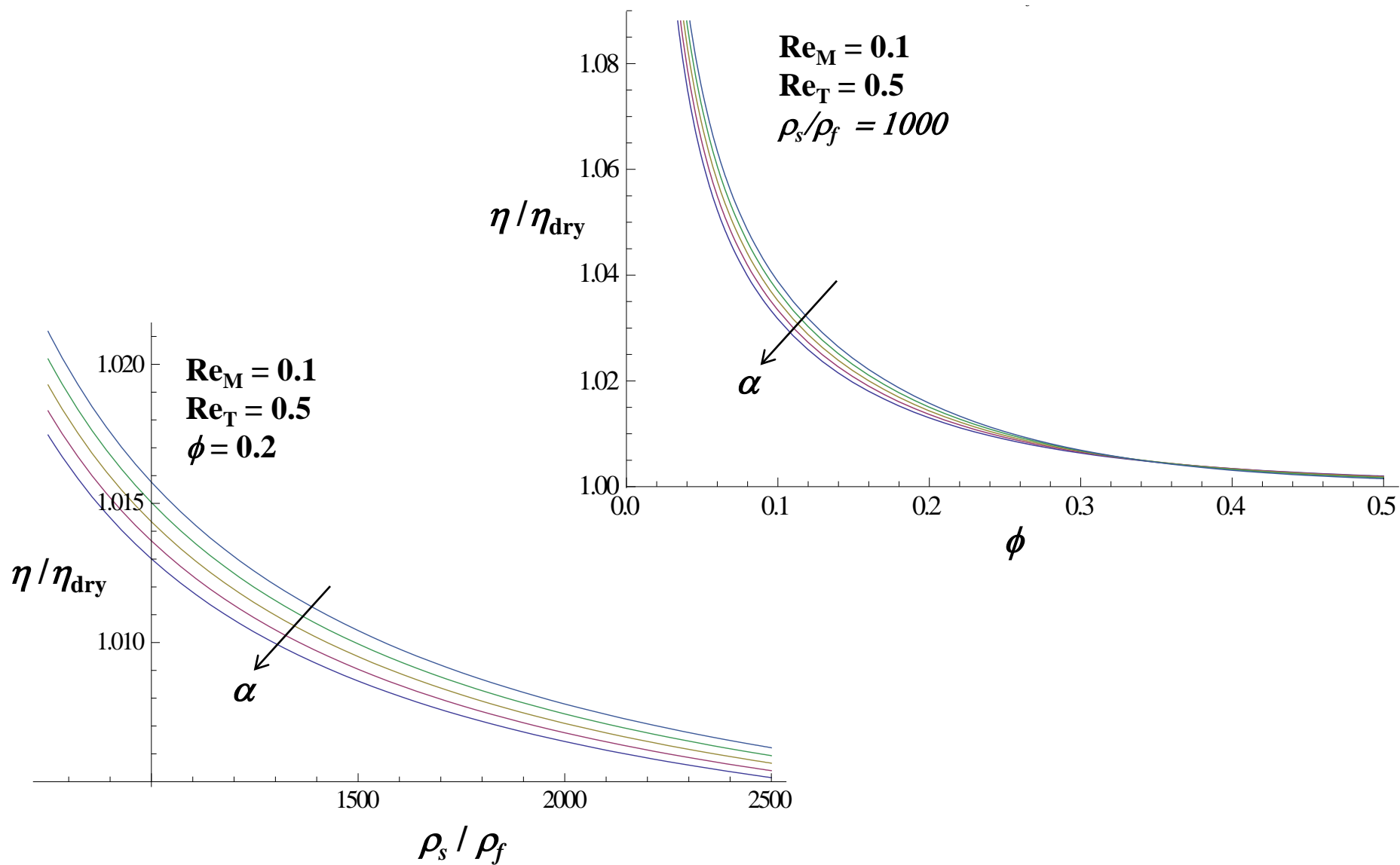
Typical Ranges in CFB (Circulating Fluidized Bed) riser

- ϕ 0.01 – 0.5
- ρ_s/ρ_f 800 – 250 \longrightarrow *high St*
- Re_m 0.1 – 50 \longrightarrow *low – moderate Re* \longrightarrow *only low $Re_m = 0.1 – 1$ considered here*
- Re_T 0.5 – 5

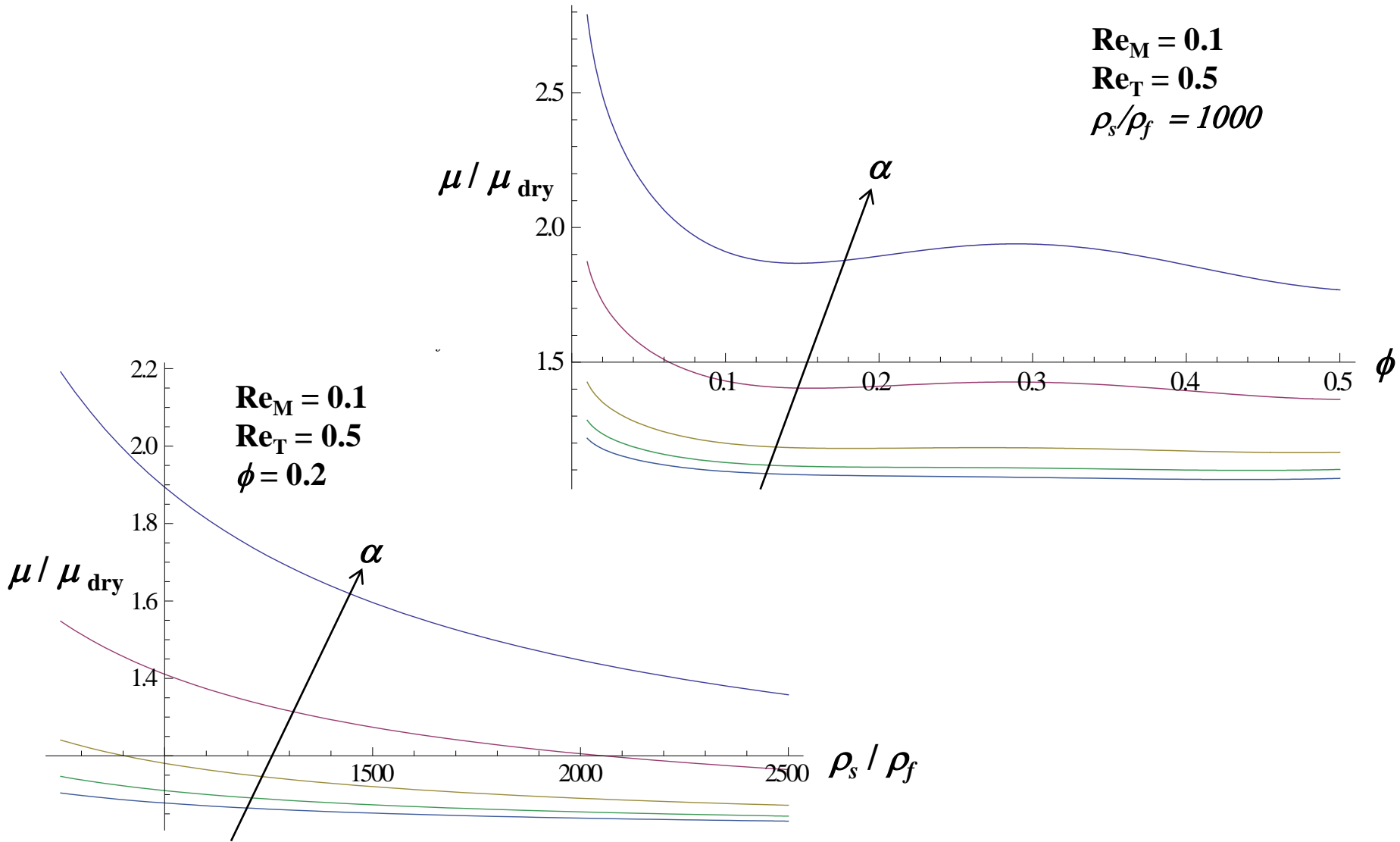
Summary of Results

- *Negligible* gas-phase influence
 - Cooling rate ($\zeta_{(0)}$, ζ_U)
 - Bulk viscosity (λ)
 - Conductivity (κ)
- *Non-negligible* gas-phase influence
 - Shear viscosity (η)
 - Dufour coefficient (μ)

Shear Viscosity



Dufour Coefficient



Summary

IBM-based model for instantaneous fluid acceleration has been incorporated into Enskog equation, and corresponding hydrodynamic description derived

- Additional source/sink in momentum and granular energy balances
- Modification of constitutive closures
 - For limiting case of $Re_m \ll 1$ and $St \gg 1$: non-negligible gas-phase influence on shear viscosity and Dufour coefficient
- Framework extendible to non-limiting cases once IBM coefficients are extracted (coming soon...)