APPLICATION OF A FILTERED EULER-LAGRANGE FORMALISM TO LARGE-SCALE SIMULATIONS OF DILUTE AND DENSE FLUID-PARTICLE FLOWS

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Multiphase flows in energy applications

- Multi-physics and multi-scale problem
 - Highly turbulent
 - Triple-phase
 - Complex geometries
 - Chemically reacting (including heat transfer + phase change)
- Computational Thermo-Fluids Laboratory led by Dr. Olivier

Desjardins

- <u>http://ctflab.mae.cornell.edu</u>
- Multi-scale and multi-physics problems
- -Plassiterfyqearatiencomputing mmersed boundaries for modeling complex geometries





Chemically reacting flows (Dr. Pepiot)



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Multi-scale issue in turbulent particle-laden flows



Macroscale

- Large number of particles $\mathcal{O}(10)$
- Length scales:



Mesoscale

- Clustering
- Bubbling
- Particle size segregation
- Turbulence modulation

Microscale

- Wakes
- Particle collisions
- Phase change





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Outline

- Filtered Euler-Lagrange framework
 - Mathematical formulation
 - Numerical implementation
- Application to dense particle-laden flows
 - Gas-solid fluidized beds
 - Liquid-solid slurries
- Application to dilute particle-laden flows
 - Turbulent channel
 - Moderately-dilute riser
- Summary & conclusions



Ingredients for developing predictive multiphase tools¹





Cornell University Computational Thermo-Fluids Laboratory 1. S. Subramaniam, Lagrangian-Eulerian methods for multiphase flows, *IJMF*, (2013)

Mathematical formulation

First-principle equations

• Gas phase: Variable-density low-Mach Navier-Stokes equations

$$\frac{\partial \rho_f}{\partial t} + \nabla \cdot (\rho_f \boldsymbol{u}_f) = 0$$

$$\frac{\partial}{\partial t} \left(\rho_f \boldsymbol{u}_f \right) + \nabla \cdot \left(\rho_f \boldsymbol{u}_f \otimes \boldsymbol{u}_f \right) = \nabla \cdot \boldsymbol{\tau} + \rho_f \boldsymbol{g}$$

$$oldsymbol{ au} = -poldsymbol{\mathcal{I}} + \mu \left[
abla oldsymbol{u}_f +
abla oldsymbol{u}_f^{\mathsf{T}} - rac{2}{3} \left(
abla \cdot oldsymbol{u}_f
ight) oldsymbol{\mathcal{I}}
ight]$$

• Particles: Newton's second law of motion

$$\begin{split} m_p \frac{d\boldsymbol{u}_p}{dt} &= \int_{\mathcal{S}_p} \boldsymbol{\tau} \cdot \boldsymbol{n} \ dS + \boldsymbol{F}_p^{\text{col}} + m_p \boldsymbol{g} \\ I_p \frac{d\boldsymbol{\omega}_p}{dt} &= \int_{\mathcal{S}_p} \frac{d_p}{2} \boldsymbol{n} \times (\boldsymbol{\tau} \cdot \boldsymbol{n}) \ d\boldsymbol{y} + \sum_j \frac{d_p}{2} \boldsymbol{n} \times \boldsymbol{f}_{t,j \to p}^{\text{col}} \end{split}$$

- Boundary conditions: no-slip and no-penetration at surface of particle
- Collision force: contact mechanics



Mathematical formulation

Volume-filtered description²

- Objective: formulate equations for particle-laden flows that allow $\Delta x \gg d_p$
- Introduce local volume filter based on convolution product with kernelg(r)
 - $\,\delta_f \gg d_p$: enabling the use of microscale models

– $\delta_f \ll \mathcal{L}_{meso}$: mesoscale structures are fully resolved



• Allows to define filtered variable \overline{a} from point variable

$$arepsilon_f \overline{oldsymbol{a}}\left(oldsymbol{x},t
ight) = \int_{\mathcal{V}_f} oldsymbol{a}\left(oldsymbol{y},t
ight) g(|oldsymbol{x}-oldsymbol{y}|) doldsymbol{y} \qquad oldsymbol{a} = \overline{oldsymbol{a}}+oldsymbol{a}')$$



2. Capecelatro & Desjardins, An Euler-Lagrange strategy for simulating particle-laden flows, *JCP*, (2012) Computational Thermo-Fluids 3. 3. T. Anderson, R. Jackson, Fluid mechanical description of fluidized beds, (1967) Laboratory 7/30

Mathematical formulation

Volume-filtering the Navier-Stokes equations

• Continuity

$$\frac{\partial}{\partial t} \left(\varepsilon_f \rho_f \right) + \nabla \cdot \left(\varepsilon_f \rho_f \overline{\boldsymbol{u}_f} \right) = 0$$

• Momentum

$$\frac{\partial}{\partial t} \left(\varepsilon_f \rho_f \overline{\boldsymbol{u}_f} \right) + \nabla \cdot \left(\varepsilon_f \rho_f \overline{\boldsymbol{u}_f} \otimes \overline{\boldsymbol{u}_f} \right) = \nabla \cdot \left(\overline{\boldsymbol{\tau}} - \boldsymbol{R}_u \right) + \varepsilon_f \rho_f \boldsymbol{g} - \boldsymbol{F}^{\text{inter}}$$
$$\overline{\boldsymbol{\tau}} = -\overline{p} \boldsymbol{\mathcal{I}} + \mu \left[\nabla \overline{\boldsymbol{u}_f} + \overline{\boldsymbol{u}_f}^{\mathsf{T}} - \frac{2}{3} \left(\nabla \cdot \overline{\boldsymbol{u}_f} \right) \boldsymbol{\mathcal{I}} \right] + \boldsymbol{R}_{\mu}$$

• Interphase exchange



Ingredients for developing predictive multiphase tools¹





Cornell University Computational Thermo-Fluids Laboratory 1. S. Subramaniam, Lagrangian-Eulerian methods for multiphase flows, *IJMF*, (2013)

Consistent framework from point-particle to full DNS

- Accurate solution of the equations requires $\Delta x \ll \delta_f$
- Model closures depend on δ_f
- What is the appropriate choice for δ_f ?

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$$\int_{\mathcal{V}_{f}} \nabla \cdot \tau g(|\mathbf{x} - \mathbf{y}|) d\mathbf{y} = \nabla \cdot (\varepsilon_{f} \overline{\tau}) - \sum_{p=1}^{n_{p}} \int_{\mathcal{S}_{p}} \mathbf{n} \cdot \tau g(|\mathbf{x} - \mathbf{y}|) d\mathbf{y}$$

$$0.01 \quad 0.1 \quad 1 \quad 10 \quad 100$$

$$\Delta x/d_{p} \xrightarrow{0.01 \quad 0.1 \quad 1 \quad 10 \quad 100} \quad \mathbf{full}$$
Fully-resolved DNS Point-particle
$$\mathbf{V} = \mathbf{V} =$$

Consistent framework from point-particle to full DNS

- Accurate solution of the equations requires $\Delta x \ll \delta_f$
- Model closures depend on δ_f

aboratory

• What is the appropriate choice for δ_f ?

$$\int_{V_f} \nabla \cdot \tau g(|x-y|) dy = \nabla \cdot (\varepsilon_f \tau) - \sum_{p=1}^{n_p} \int_{S} (2p_1 + p_2) dy = \nabla \cdot (\varepsilon_f \tau) - \sum_{p=1}^{n_p} \int_{S} (2p_1 + p_2) dy = \nabla \cdot (\varepsilon_f \tau) - \sum_{p=1}^{n_p} \int_{S} (2p_1 + p_2) dy = \nabla \cdot (\varepsilon_f \tau) - \sum_{p=1}^{n_p} \int_{S} (2p_1 + p_2) dy = \nabla \cdot (\varepsilon_f \tau) - \sum_{p=1}^{n_p} \int_{S} (2p_1 + p_2) dy = \nabla \cdot (\varepsilon_f \tau) - \sum_{p=1}^{n_p} \int_{S} (2p_1 + p_2) dy = \nabla \cdot (\varepsilon_f \tau) - \sum_{p=1}^{n_p} \int_{S} (2p_1 + p_2) dy = \nabla \cdot (\varepsilon_f \tau) - \sum_{p=1}^{n_p} \int_{S} (2p_1 + p_2) dy = \nabla \cdot (\varepsilon_f \tau) - \sum_{p=1}^{n_p} \int_{S} (2p_1 + p_2) dy = \nabla \cdot (\varepsilon_f \tau) - \sum_{p=1}^{n_p} \int_{S} (2p_1 + p_2) dy = \nabla \cdot (\varepsilon_f \tau) - \sum_{p=1}^{n_p} \int_{S} (2p_1 + p_2) dy = \nabla \cdot (\varepsilon_f \tau) - \sum_{p=1}^{n_p} \int_{S} (2p_1 + p_2) dy = \nabla \cdot (\varepsilon_f \tau) - \sum_{p=1}^{n_p} \int_{S} (2p_1 + p_2) dy = \nabla \cdot (\varepsilon_f \tau) - \sum_{p=1}^{n_p} \int_{S} (2p_1 + p_2) dy = \nabla \cdot (\varepsilon_f \tau) - \sum_{p=1}^{n_p} \int_{S} (2p_1 + p_2) dy = \nabla \cdot (\varepsilon_f \tau) - \sum_{p=1}^{n_p} \int_{S} (2p_1 + p_2) dy = \nabla \cdot (\varepsilon_f \tau) - \sum_{p=1}^{n_p} \int_{S} (2p_1 + p_2) dy = \nabla \cdot (\varepsilon_f \tau) - \sum_{p=1}^{n_p} \int_{S} (2p_1 + p_2) dy = \nabla \cdot (\varepsilon_f \tau) - \sum_{p=1}^{n_p} \int_{S} (2p_1 + p_2) dy = \nabla \cdot (\varepsilon_f \tau) - \sum_{p=1}^{n_p} \int_{S} (2p_1 + p_2) dy = \nabla \cdot (\varepsilon_f \tau) - \sum_{p=1}^{n_p} \int_{S} (2p_1 + p_2) dy = \nabla \cdot (\varepsilon_f \tau) - \sum_{p=1}^{n_p} \int_{S} (2p_1 + p_2) dy = \nabla \cdot (\varepsilon_f \tau) - \sum_{p=1}^{n_p} \int_{S} (2p_1 + p_2) dy = \nabla \cdot (\varepsilon_f \tau) - \sum_{p=1}^{n_p} \int_{S} (2p_1 + p_2) dy = \nabla \cdot (\varepsilon_f \tau) - \sum_{p=1}^{n_p} \int_{S} (2p_1 + p_2) dy = \nabla \cdot (\varepsilon_f \tau) - \sum_{p=1}^{n_p} \int_{S} (2p_1 + p_2) dy = \nabla \cdot (\varepsilon_f \tau) - \sum_{p=1}^{n_p} \int_{S} (2p_1 + p_2) dy = \nabla \cdot (\varepsilon_f \tau) - \sum_{p=1}^{n_p} \int_{S} (2p_1 + p_2) dy = \nabla \cdot (2p_1$$

Numerical implementation

Filter discretization

Laboratory

- Direct implementation becomes too expensive
- Filter based on the convolution of mollification and Laplacian smoothing²
 - 1. Mollification: transfer particle data to neighboring cells
 - 2. Diffusion: smooth data with specified width
- Fully conservative, implicit treatment
- Special care is needed at the walls



Numerical implementation



Ingredients for developing predictive multiphase tools¹





Cornell University Computational Thermo-Fluids Laboratory 1. S. Subramaniam, Lagrangian-Eulerian methods for multiphase flows, *IJMF*, (2013)

Ingredients for developing predictive multiphase tools¹





Cornell University Computational Thermo-Fluids Laboratory 1. S. Subramaniam, Lagrangian-Eulerian methods for multiphase flows, *IJMF*, (2013)

Numerical implementation

Computational platform

NGA⁴

- Arbitrarily high-order multi-physics DNS/LES code
- Conservation of mass, momentum, and kinetic energy
- Highly scalable





4. O. Desjardins, G. Cornell University Computational Thermo-Fluids Laboratory

4. O. Desjardins, G. Blanquart, G. Balarac, H. Pitsch, High order conservative finite difference scheme for variable density low Mach number turbulent flows, *JCP* (2008)

Ingredients for developing predictive multiphase tools¹





Cornell University Computational Thermo-Fluids Laboratory 1. S. Subramaniam, Lagrangian-Eulerian methods for multiphase flows, *IJMF*, (2013)

Application to dense gas-solid flows²



Application to liquid-solid slurries⁵

- Compared with experiments by Roco & Balakrishnam (1985)
- Two cases simulated
 - Re=85,000 (above critical deposition velocity)
 - Re=42,660 (below critical deposition velocity)
- Force liquid mass flow rate in a periodic pipe
- 768 x 156 x 156 mesh
- 19 M polydisperse particles





Cornell University Computational Thermo-Fluids Laboratory 5. Capecelatro & Desjardins, Eulerian-Lagrangian modeling of turbulent liquid-solid slurries in horizontal pipes, *IJMF* (2013)

Application to liquid-solid slurries

Computational Thermo-Fluids Lab

Time = 0.0000





Application to liquid-solid slurries

Laboratory



Application to liquid-solid slurries



Simulation configuration



- δ = 2 cm
- Re₂₀ = 13,850
- Bulk velocity = 9.2 m/s
- Centerline velocity, $U_{cl} = 10.5$
- Particle diameter = 150 µm
- Particle density = 2,500 kg/m³
- St = 50



	Re_{τ}	ϕ	Wall boundary condition
Paris & Eaton (2001)	644	0.2	Rough
Benson & Eaton (2003)	617	0.15	Smooth
NGA	630	0.15	Smooth

Single-phase channel flow





Particle-laden channel





Turbulence modulation & preferential concentration due to non-uniform interphase coupling



Particle number density



Particle-laden channel





Application to risers

• Dimensional analysis

$${\rm Fr} = U/\sqrt{gd_p} \qquad {\rm Ar} = \rho_s \rho_f d_p^3 g/\mu^2 \quad D/d_p$$

- Experimental observations (Noymer & Glicksman, 2000)
 - Clusters fall very close to the walls (~100 μm)
 - Clusters located within hydrodynamic boundary layer
 - Cluster fall velocity independent of inflow conditions

$$\frac{u_{cl}}{u_{mf}} = \frac{1000}{\sqrt{\text{Ar}}} \quad u_{mf} = 0.00075 \frac{\rho_s g d_p^2}{\mu}$$

$$u_{cl} = 0.75 \sqrt{\frac{\rho_s}{\rho_f} g d_p}$$

- Simulation parameters
 - 3D pipe geometry (immersed boundaries)
 - Periodic in vertical direction
 - 760,000 particles
 - Mesh: 800x83x83



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	Case 1	Case 2	Case 3	Case 4	Case 5				
Ar	50	100	500	2500	12500				
D/d_p	320	150	150	150	150				
$\langle \varepsilon_p \rangle$	0.15%	1.5%	1.5%	1.5%	1.5%				
$ ho_p/ ho_f$	2500	2500	2500	2500	2500				

Simulation cases



Application to risers



Cluster fall velocity

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Conclusions

- Volume filtered formalism provides a consistent framework from point-particle to fully resolved simulations
- Pushing the validity of classical microscale models to finer meshes yields excellent results
- The proposed framework can capture a range of phenomenon including
 - Clustering
 - Bubbling
 - Segregation in particle size
 - Preferential concentration
- Looking forward
 - Study intermediate values of particle diameter to mesh size ratio ($\lesssim d_p/\Delta x \lesssim 10$)
 - Implement sharper / higher accuracy filters
 - Use this framework to provide closure for RANS modeling:

R.O. Fox, J. Capecelatro, O. Desjardins Validation of a Multiphase Turbulence Model Using Mesoscale DNS of Gravity-Driven Gas-Particle Flow. 11:50-12:10 PM

