

An Easily Implementable General Self-induced Perturbation Correction Model for a Finite-sized Particle in Two-way Coupled Euler-Lagrange Simulations

Kai Liu^{a,b}, S. Balachandar^a

^a University of Florida

^b Zhejiang University



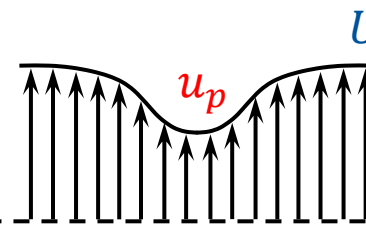
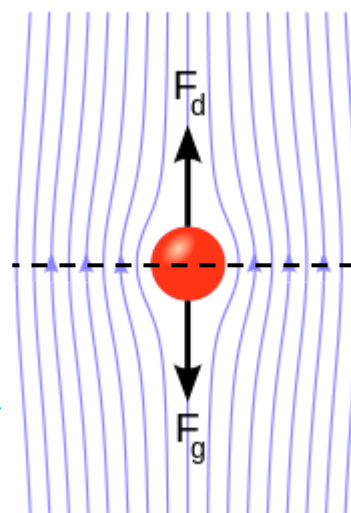
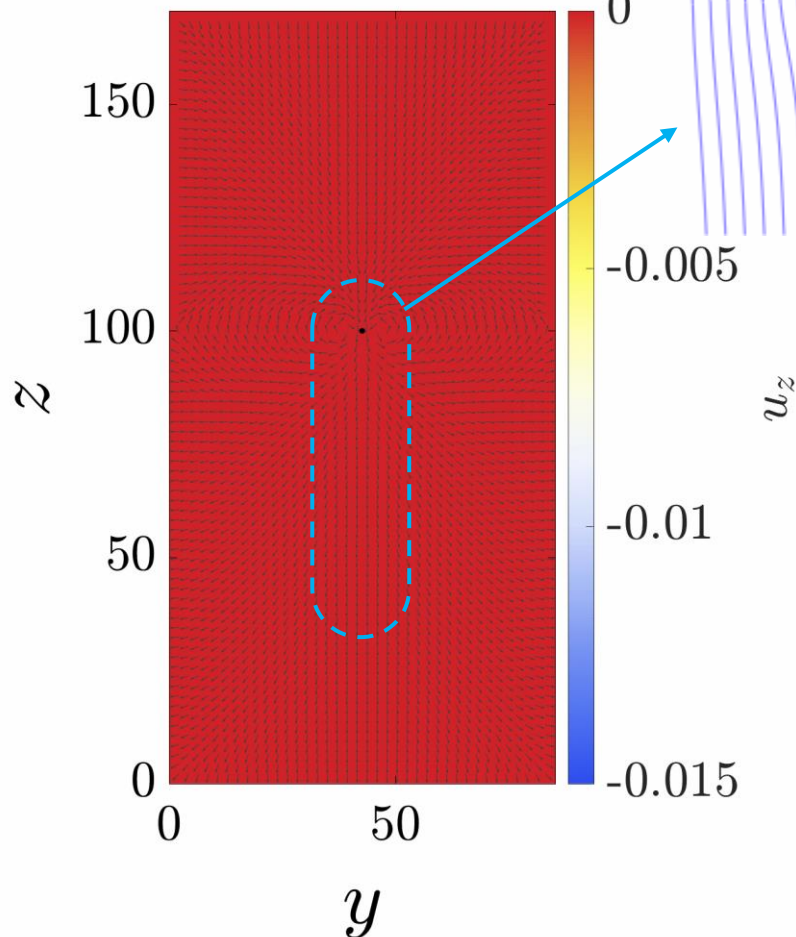
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Motivation

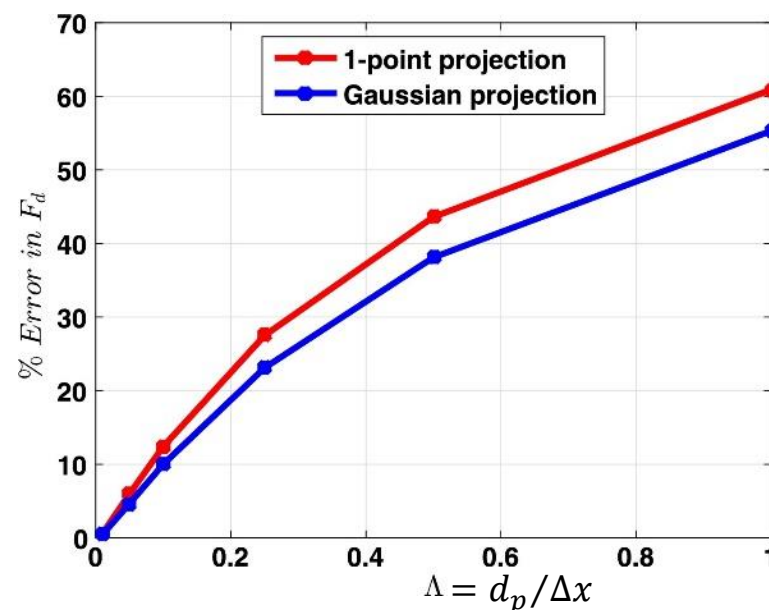
single particle settling

$Re = 260.6$ $t = 0$

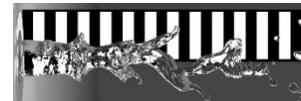


Stokes' drag:

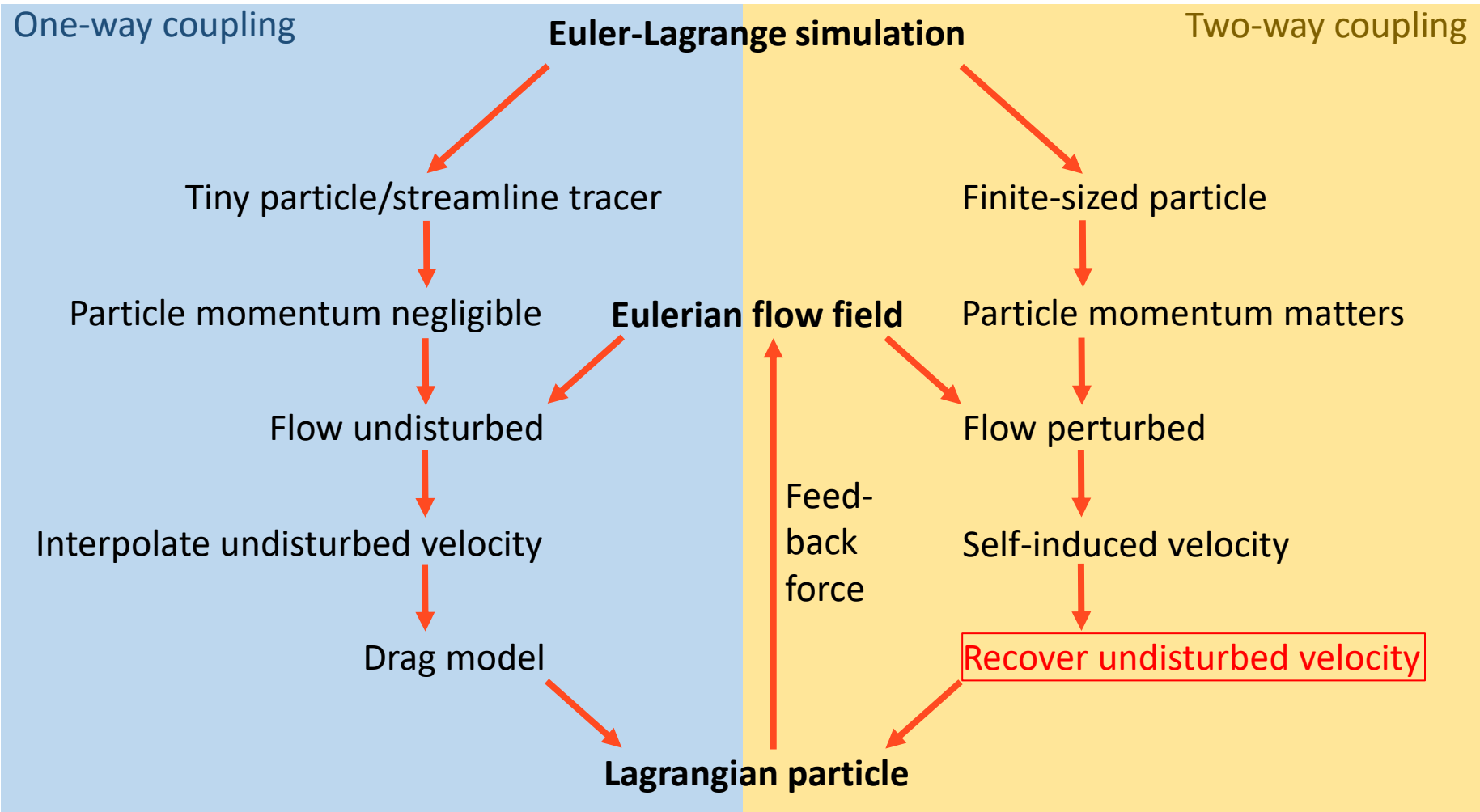
$$F_d = \begin{cases} 3\pi\mu d_p(U - v_p) & \checkmark \text{ By theory} \\ 3\pi\mu d_p(u_p - v_p) & \times \text{ In simulation} \end{cases}$$



J. A. K. Horwitz and A. Mani. Journal of Computational Physics, 2016.



Hierarchy



Our plan

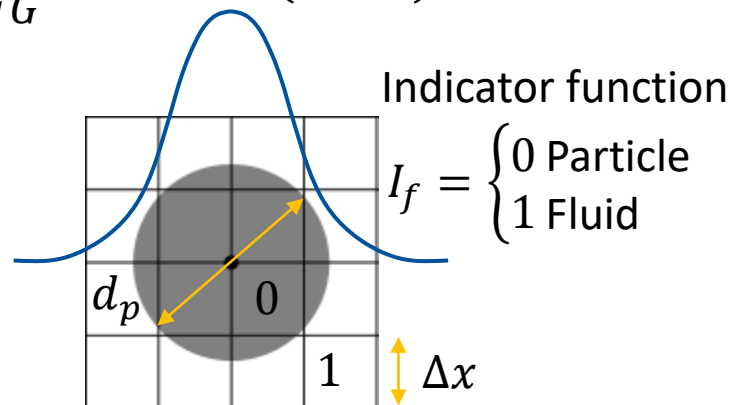
- Simple model
- Drag not in same direction of flow
- Near wall
- Compatible for common projection filters
- Compatible for common difference schemes

Oseen's equation:

$$\nabla \cdot (\phi_{c1} \mathbf{u}'_{1c}) = 0$$

$$\rho \mathbf{U} \cdot (\phi_{c1} \mathbf{u}'_{1c}) = -\nabla(\phi_{c1} p'_{1c}) + \mu \nabla^2(\phi_{c1} \mathbf{u}'_{1c}) - \mathbf{F}G$$

$$G(\mathbf{r}) = \frac{1}{(\sqrt{2\pi}\sigma)^3} \exp\left(-\frac{r^2}{2\sigma^2}\right)$$

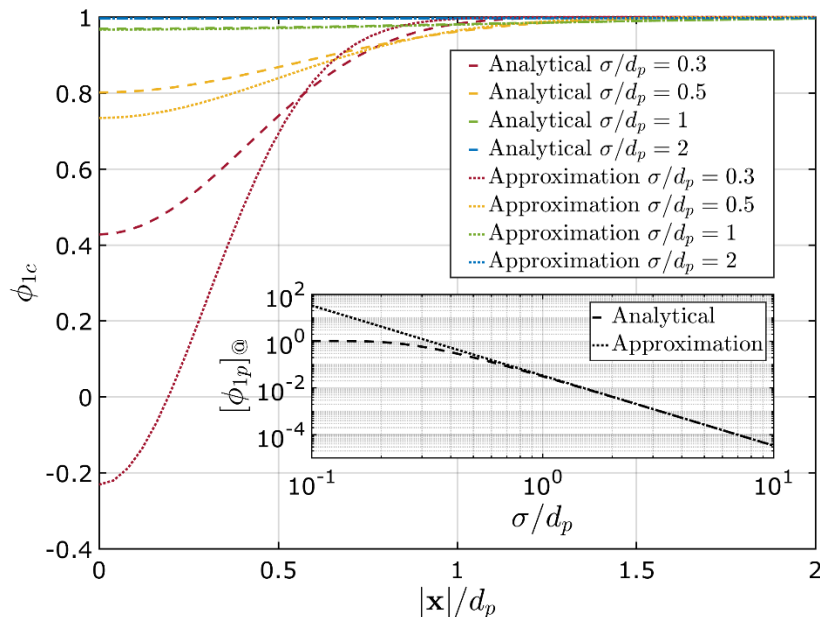


Solution:

$$\underbrace{\mathbf{U}(\mathbf{x}, t)}_{\text{Undisturbed velocity}} = \underbrace{\phi_{c1} \mathbf{u}_c(\mathbf{x}, t)}_{\text{EL particle Center velocity}} - \underbrace{\phi_{c1} \mathbf{u}'_{1c}(\mathbf{x}, t)}_{\text{Self-induced perturbation velocity}}$$



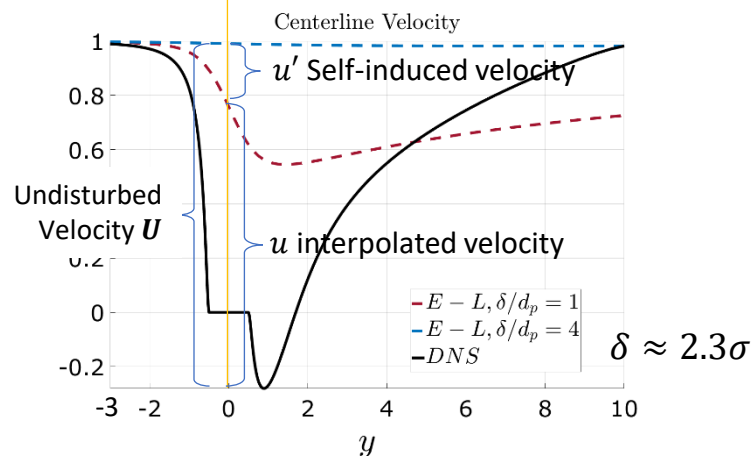
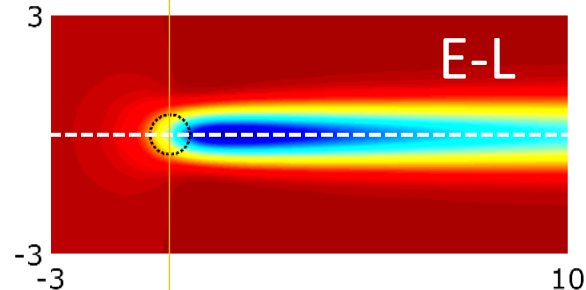
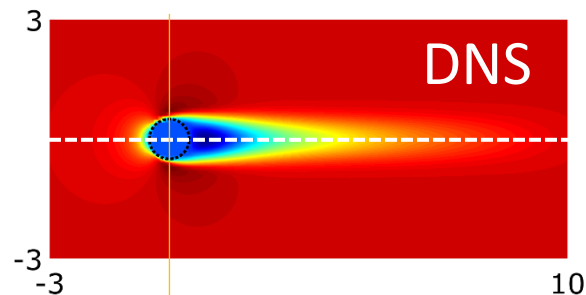
Filter effect



$$\phi_{1c} = \begin{cases} 1 - \operatorname{erf}\left(\frac{d_p}{2\sqrt{2}\sigma}\right) + \frac{d_p}{\sqrt{2\pi}\sigma} \exp\left(-\frac{d_p^2}{8\sigma^2}\right) & \text{Analytical} \\ 1 - \frac{\pi d_p^3}{6} \frac{1}{(\sqrt{2\pi}\sigma)^3} \exp\left(-\frac{0^2}{2\sigma^2}\right) & \text{Approximation} \end{cases}$$

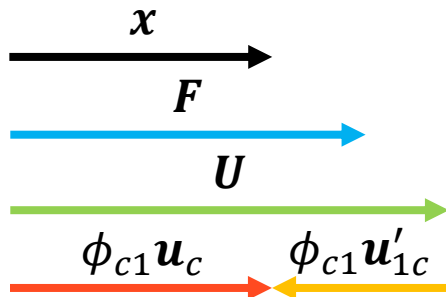
- Recommend using the analytical form
- $\sigma/d_p \rightarrow$ fluid volume fraction
- $\sigma/d_p \rightarrow$ self-induced velocity

$$U = 1, d_p = 1, Re_p = 170$$

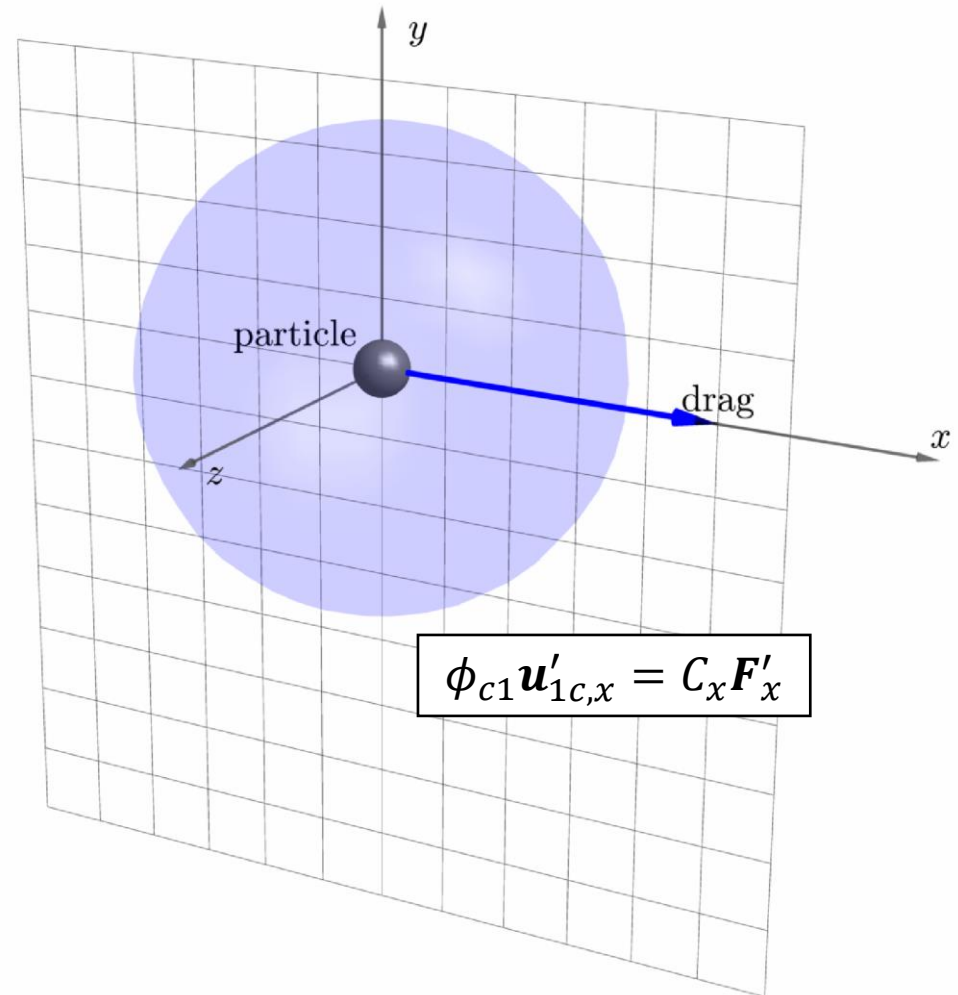
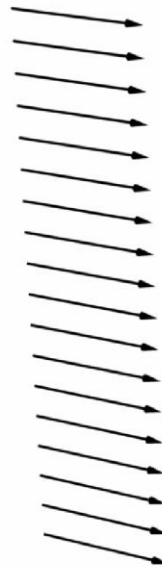


Scenario 1: x-dir scalar

- Drag is parallel to flow
- self-induced velocity in x-dir
- No y,z perturbation @ $x = 0$



far-field flow



$$\phi_{c1} u'_{1c,x} = C_x F'_x$$

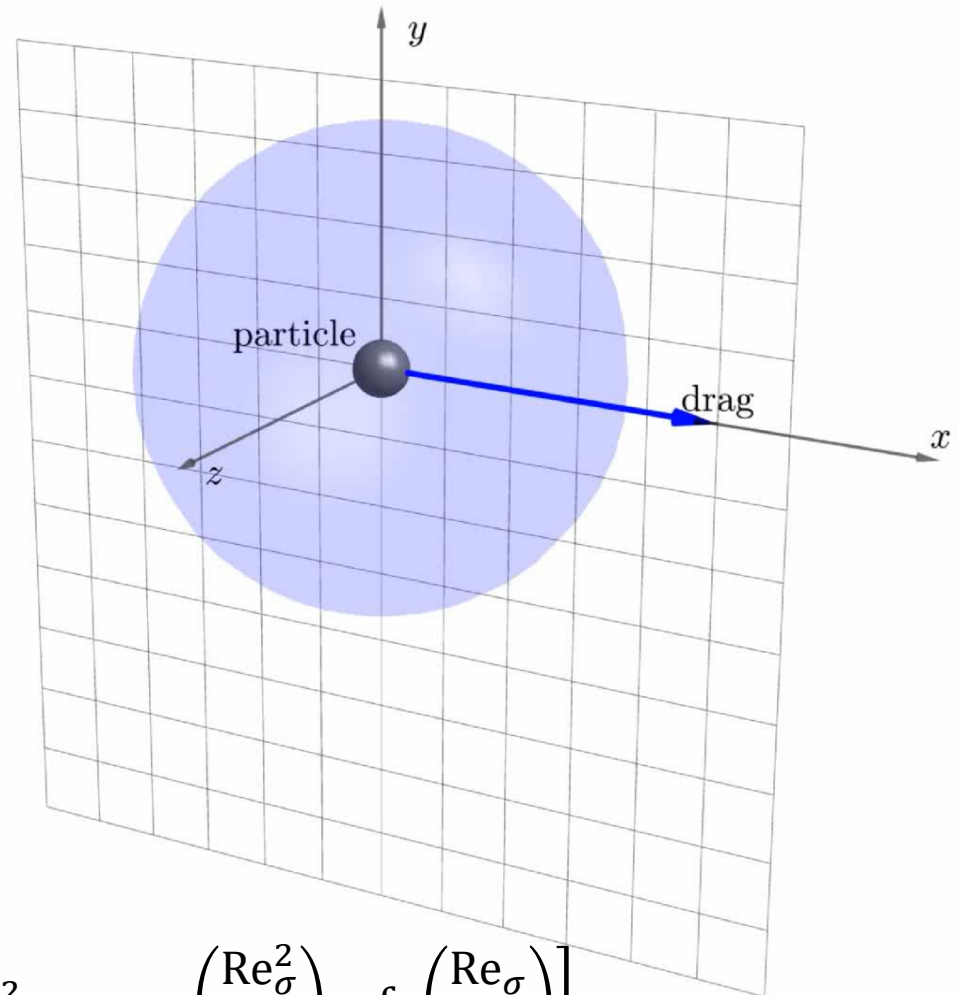
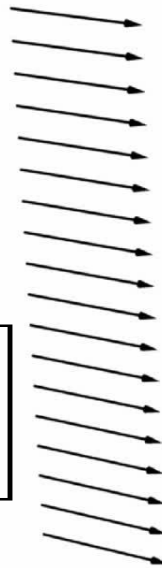


Scenario 2: 3D vectorial

- Drag is not parallel to flow
- Due to not only \mathbf{F}'_{qs}
- Diagonal relation

$$\begin{bmatrix} \phi_{c1} \mathbf{u}'_{1c,x} \\ \phi_{c1} \mathbf{u}'_{1c,y} \\ \phi_{c1} \mathbf{u}'_{1c,z} \end{bmatrix} = \begin{bmatrix} M_x & 0 & 0 \\ 0 & M_y & 0 \\ 0 & 0 & M_z \end{bmatrix} \begin{bmatrix} \mathbf{F}'_x \\ \mathbf{F}'_y \\ \mathbf{F}'_z \end{bmatrix}$$

far-field flow

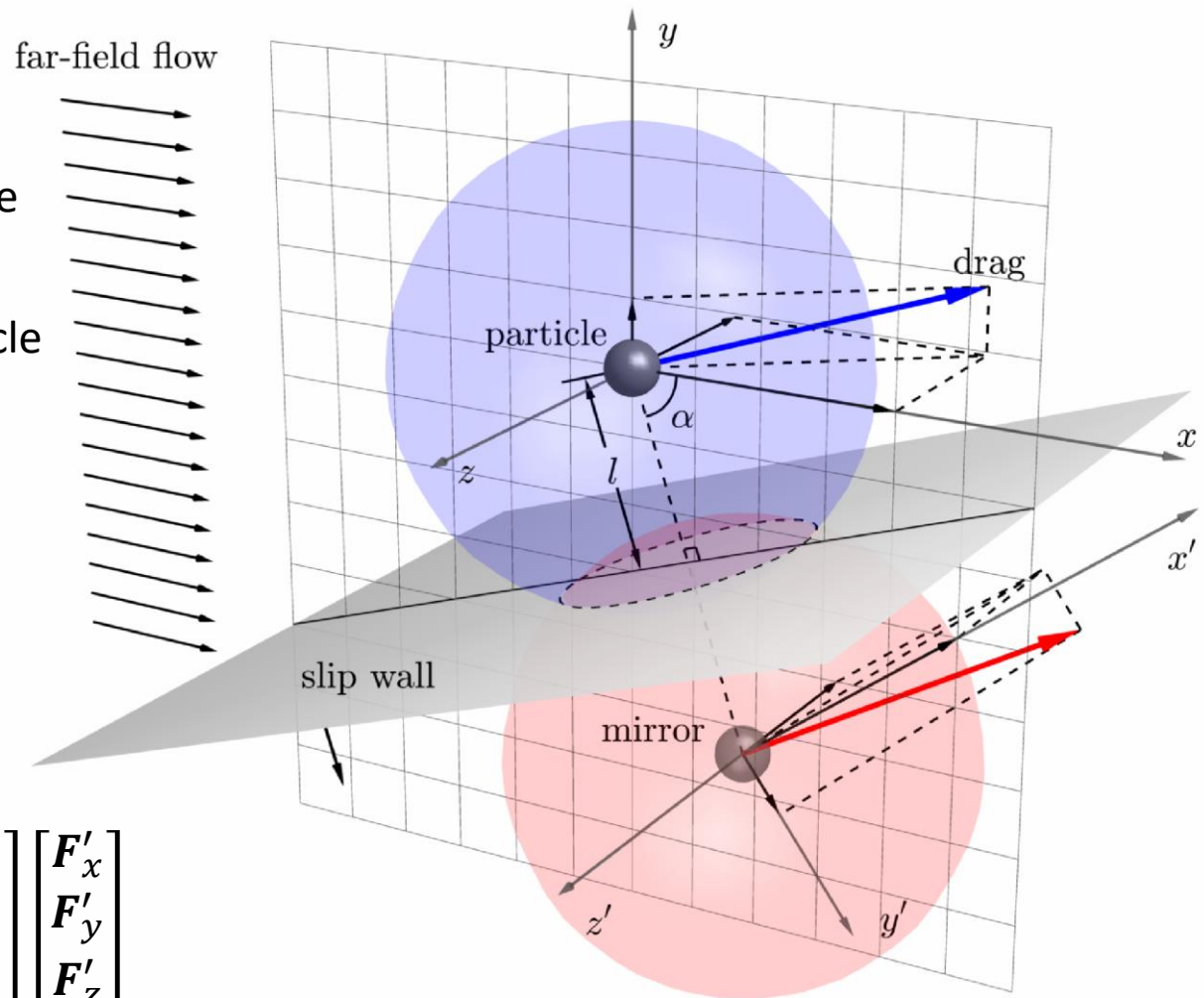


$$\text{Re}_\sigma = \frac{|\bar{\mathbf{U}}^s(\mathbf{x} = \mathbf{0}, t) - \mathbf{V}| \sigma}{\nu}$$

$$M_x(\text{Re}_\sigma) = -\frac{1}{2\pi^2 \text{Re}_\sigma^2} \left[\pi - \sqrt{2\pi} \text{Re}_\sigma + \frac{\pi}{2} \text{Re}_\sigma^2 - \pi \exp\left(\frac{\text{Re}_\sigma^2}{2}\right) \text{erfc}\left(\frac{\text{Re}_\sigma}{\sqrt{2}}\right) \right]$$

Scenario 3: near-wall

- Free-slip wall nearby
- Projection point in x-y plane
- Modelled by a mirror particle
- Symmetric by the wall
- Drag symmetric by the wall
- Not only diagonal relation



$$\begin{bmatrix} \phi_{c1} \mathbf{u}'_{1c,x} \\ \phi_{c1} \mathbf{u}'_{1c,y} \\ \phi_{c1} \mathbf{u}'_{1c,z} \end{bmatrix} = \begin{bmatrix} M_{xx} & M_{xy} & 0 \\ M_{yx} & M_{yy} & 0 \\ 0 & 0 & M_{zz} \end{bmatrix} \begin{bmatrix} \mathbf{F}'_x \\ \mathbf{F}'_y \\ \mathbf{F}'_z \end{bmatrix}$$

Correction model

standard drag

$$\mathbf{F}'_{qs} = 3\pi\mu d_p [\overline{\phi_{c1} \mathbf{u}_c^s} - \mathbf{V}] \Phi \quad \text{Interpolated relative flow velocity } \overline{\phi_{c1} \mathbf{u}_c^s} - \mathbf{V}$$

Renew
 $\bar{\mathbf{U}}^s - \mathbf{V}$
and
Iterate

Coordinates rotation $[\mathbf{M}][\overline{\phi_{c1} \mathbf{u}_c^s} - \mathbf{V}]$

Compute $\text{Re}_\sigma, \text{Re}_p, \Phi, [\mathbf{B}], [\mathbf{C}]$ using $[\mathbf{M}][\bar{\mathbf{U}}^s - \mathbf{V}]$

$$\bar{\mathbf{U}}^s - \mathbf{V} = [\mathbf{M}]^T \left\{ [\mathbf{I}] - \frac{d_p}{\sigma} \frac{\Phi}{\sqrt{2\pi}} [\mathbf{B}]([\mathbf{I}] + [\mathbf{C}]) \right\}^{-1} [\mathbf{M}](\overline{\phi_{c1} \mathbf{u}_c^s} - \mathbf{V})$$

Converge

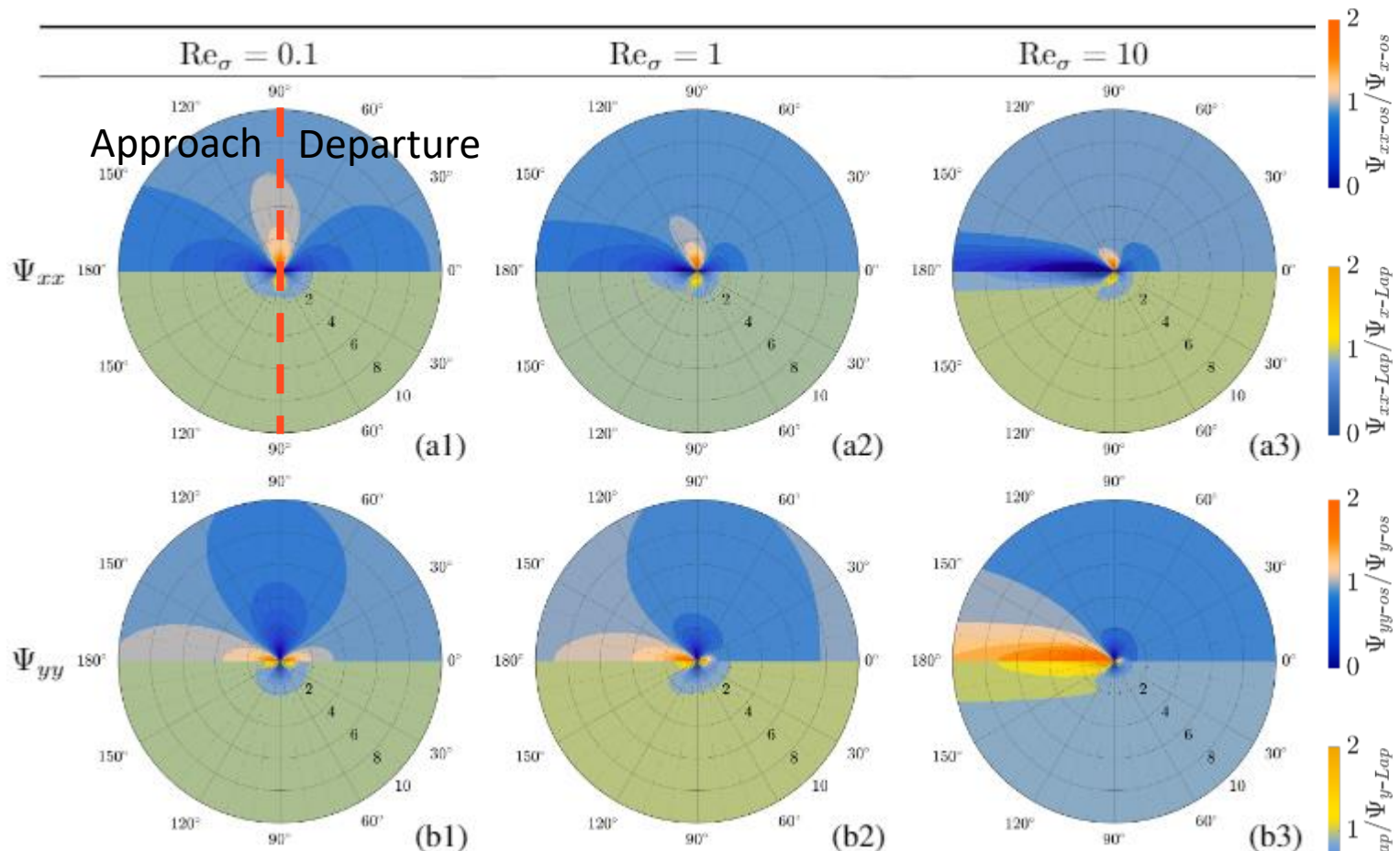
True drag $\mathbf{F}'_{qs} = 3\pi\mu d_p [\bar{\mathbf{U}}^s - \mathbf{V}] \Phi$

$$[\mathbf{B}] = \begin{bmatrix} \Psi_{xx-os} & \Psi_{xy-os} & 0 \\ \Psi_{yx-os} & \Psi_{yy-os} & 0 \\ 0 & 0 & \Psi_{zz-os} \end{bmatrix} - \frac{d_p^2}{24\sigma^2} \begin{bmatrix} \Psi_{xx-Lap} & \Psi_{xy-Lap} & 0 \\ \Psi_{yx-Lap} & \Psi_{yy-Lap} & 0 \\ 0 & 0 & \Psi_{zz-Lap} \end{bmatrix}$$

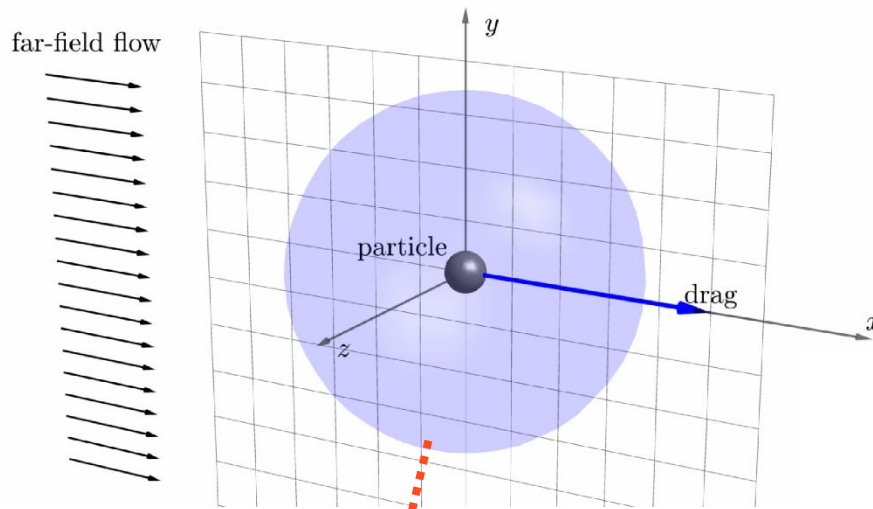
$$[\mathbf{C}] = [\mathbf{M}] \left[\frac{\mathbf{F}'_{ot}}{\mathbf{F}'_{qs}} \right] \left[\frac{\overline{\phi_{c1} \mathbf{u}_c^s} - \mathbf{V}}{\bar{\mathbf{U}}^s - \mathbf{V}} \right] [\mathbf{M}]^T$$

Ψ components

Obtained totally 10 components of $\Psi \left(\text{Re}_\sigma, \frac{l}{\sigma}, \alpha \right)$ with explicit asymptotic functions for low Re_σ



Different filters



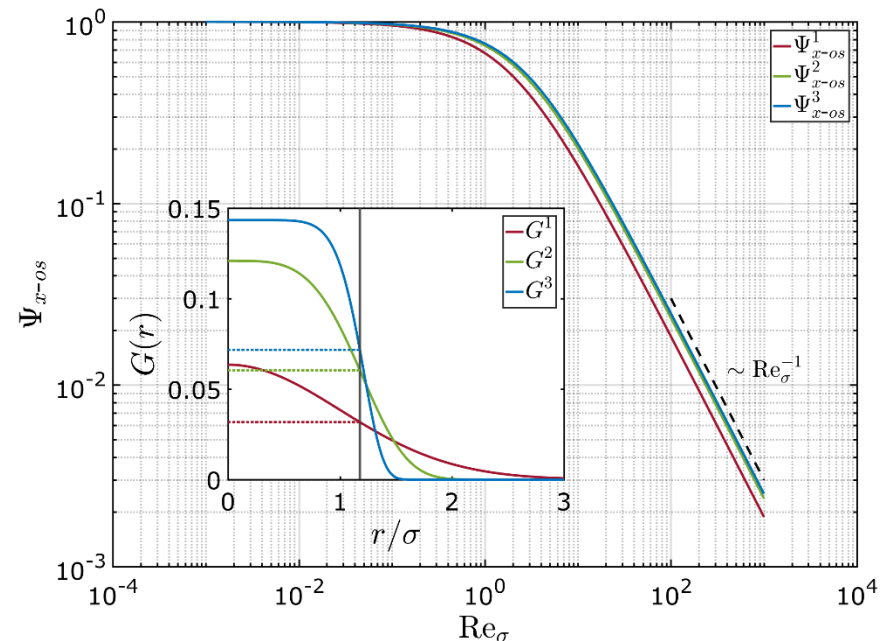
$$G^1(r) = \frac{1}{(2\pi)^{3/2}\sigma^3} \exp\left(-\frac{r^2}{2\sigma^2}\right)$$

$$G^2(r) = \frac{1}{\pi\Gamma\left(\frac{3}{4}\right)(\ln 16)^{3/4}\sigma^3} \exp\left[-\frac{r^4}{(\ln 16)\sigma^4}\right]$$

$$G^3(r) = \frac{2^{5/8}}{\pi\Gamma\left(\frac{3}{8}\right)(\ln 4)^{9/8}\sigma^3} \exp\left[-\frac{r^8}{2(\ln 4)^3\sigma^8}\right]$$

Scalar case for example, different choice of filter does not change general solution

$\overline{\phi_{c1} \mathbf{u}'_c}^s = C \text{Re}_\sigma \mathbf{F}'_x \Psi_{x-os}$
but slightly changes the specific values of C, Ψ_{x-os}



Differentiation schemes

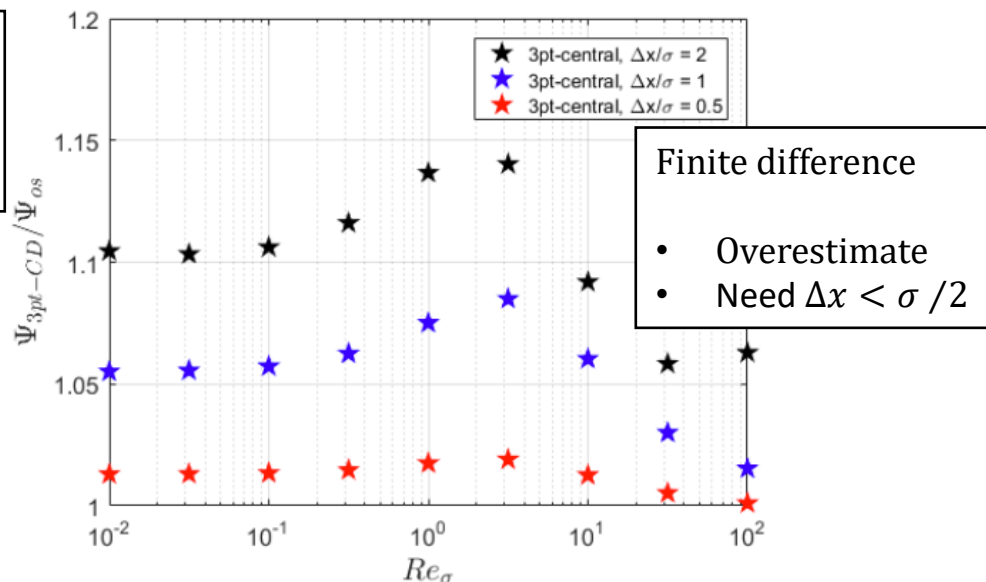
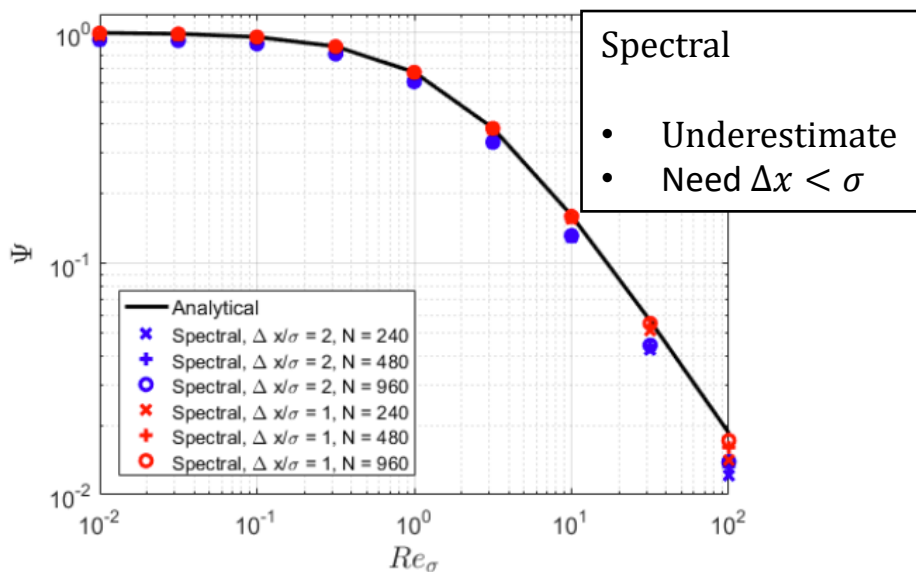
Self-induced velocity derived by Fourier Transform $\widehat{\phi_{c1} \mathbf{u}'_c} = - \frac{F'_x \text{Re}_\sigma \left(1 - i \frac{\text{Re}_\sigma \tilde{k}_x}{|\tilde{k}|^2} \right)}{|\tilde{k}|^2 \left(1 + \frac{\text{Re}_\sigma^2 \tilde{k}^2}{|\tilde{k}|^4} \right)} \left(1 - \frac{\tilde{k}_x^2}{\tilde{k}^2} \right) \hat{G}$

Wave number in spectral method

$$\tilde{k} = k \text{ and } \tilde{k}^2 = k^2 \text{ for } k = 0, \pm k_{\min}, \pm 2k_{\min}, \dots, \pm \frac{N}{2} k_{\min}$$

Wave number in second-order finite difference scheme

$$\tilde{k} = \frac{1}{\Delta x} \sin(k\Delta x) \text{ and } \tilde{k}^2 = \frac{1}{\Delta x^2} [1 - \cos(k\Delta x)] \text{ for } k = 0, \pm k_{\min}, \pm 2k_{\min}, \dots, \pm \frac{N}{2} k_{\min}$$



Conclusion

Subject: provide the self-induced perturbation velocity correction model

Aim: recover the true drag force in Two-way coupled Euler-Lagrange simulation

Needs: $\Delta x < \sigma$ for spectral method, $\Delta x < \sigma / 2$ for finite difference scheme

Inputs: $\sigma, d_p, \nu, \mathbf{X}_w, (\overline{\phi_{c1} \mathbf{u}_c}^s - \mathbf{V}), \left[\frac{\mathbf{F}'_{ot}}{\mathbf{F}'_{qs}} \right]$

Outputs: corrected drag

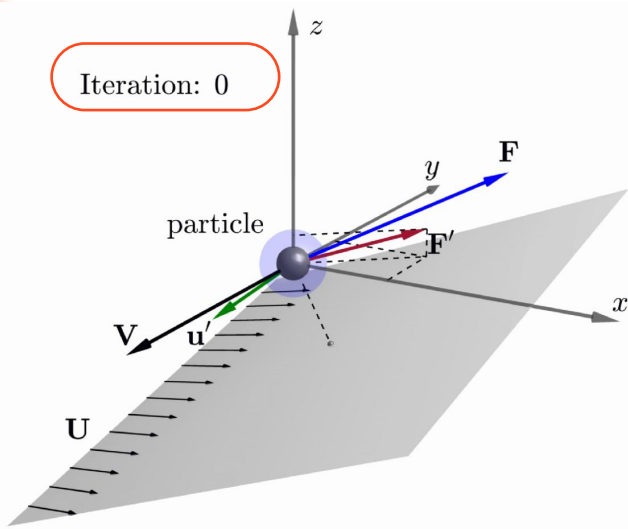
Capability:

- Arbitrary velocity vectors of flow, particle, and force
- Free-slip wall

Feature:

- Analytically derived
- Explicit form without table interpolation
- Easy to extend to different filters, DS, finite Re_σ , fast-varying, heat/mass transfer

Iteration: 0



Thank you for watching!

You can send me questions by email: kailiu@ufl.edu



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ONR MURI
Spray Control



Extension

1. Finite Reynolds number $Re_\sigma \gg 1$ correction
 - Moderate derivation
 - Small self-induced velocity, less important to correct
2. Fast transient correction
 - Derive time history term
 - Commonly not used
3. Self-induced temperature/concentration corrections
 - Scalar form
 - Adiabatic/impermeable wall

