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On the effect of particle Froude number in sub-grid modeling of gas-solid fluidized flows

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Our ultimate objective (constitutive modeling through HRS)





At left a numerical mesh of a Large Scale Simulation (LSS), highliting a grid cell over which the inside meso-scale flow is dissipated and turned homogeneous.

At right a numerical mesh of a Highly Resolved Simulation (HRS) applying fine grids thereby allowing the capture of meso-scale flow heterogeneities, also showing a space filter applied over the grids so as to recover effects of flow heterogeneities to be provided to the Large Scale Simulations as constitutive models.



Microscopic two-fluid model



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$$\frac{\partial}{\partial t} (\rho_g \phi_g) + \nabla \cdot (\rho_g \phi_g v_g) = 0$$
$$\frac{\partial}{\partial t} (\rho_s \phi_s) + \nabla \cdot (\rho_s \phi_s v_s) = 0$$
$$\frac{\partial}{\partial t} (\rho_g \phi_g v_g) + \nabla \cdot (\rho_g \phi_g v_g v_g) = -\phi_g \nabla \cdot \sigma_g - M_l + \rho_g \phi_g g$$
$$\frac{\partial}{\partial t} (\rho_s \phi_s v_s) + \nabla \cdot (\rho_s \phi_s v_s v_s) = -\nabla \cdot \sigma_s - \phi_s \nabla \cdot \sigma_g + M_l + \rho_s \phi_s g$$
$$\phi_g + \phi_s = 1$$









$$\frac{3}{2} \left[\frac{\partial}{\partial t} (\rho_s \phi_s \theta) + \nabla \cdot (\rho_s \phi_s \boldsymbol{v}_s \theta) \right] = -\boldsymbol{\sigma}_s : \nabla \boldsymbol{v}_s + \nabla \cdot (k_s \theta) + \Gamma_{slip} - J_{coll} - J_{vis}$$

$$p = \frac{|\boldsymbol{v}_g - \boldsymbol{v}_s| u_p p_g}{\mu_g}$$



Filtered parameters



$$H = 1 - \frac{\beta_{eff}}{\bar{\beta}}$$
$$\beta_{eff} = \frac{\overline{\beta(\boldsymbol{v}_g - \boldsymbol{v}_s)}}{(\tilde{\boldsymbol{v}}_g - \tilde{\boldsymbol{v}}_s)} - \frac{[\overline{\phi_s \nabla P_g} - \overline{\phi}_s \nabla \tilde{P}_g]}{(\tilde{\boldsymbol{v}}_g - \tilde{\boldsymbol{v}}_s)}$$

$$P_{fill,s} = \overline{P}_{s} - \overline{\left(\lambda_{s} + \frac{2}{3}\mu_{s}\right)}(\nabla \cdot \boldsymbol{v}_{s})$$

$$\mu_{fill,s} = \bar{\mu}_s$$



Residual parameters



$$P_{res,l} = \frac{1}{3} tr(\boldsymbol{\tau}_{l}')$$

$$\mu_{res,l} = \frac{|\boldsymbol{\tau}_{shear,l}'|}{2|\tilde{S}_{shear,l}'|}$$



HRS simulations



Simulations were performed with the microscopic two-fluid model of MFIX, for:

- low Reynolds number suspensions
- 15% domain average solid volume fraction
- particulates of 40, 75, 150, 300 μ m (Fr_{dp}= 12.21, 64.95, 286.69, 799.26)
- all periodic boundaries
- 16 x 16 cm domain
- 128 x 128 grids (1.25 x 1.25 mm grid cells)

Particulate and gas properties:

$$\label{eq:rhog} \begin{split} \rho_s &= 1500 \ \text{kg/m}^3 & e &= 0.9 \\ \rho_g &= 1.3 \ \text{kg/m}^3 & \mu_g &= 1.8 \times 10^{-5} \ \text{kg/(m s)} \end{split}$$





HRS snapshot (grayscale plot of solid volume fraction)

Filtering HRS data



- A square filter box is defined embracing a number of HRS grid cells, which is made to sweep all over the domain while calculating inside averages (i.e. filtered data).
- Filtered data are classified while statistically averaged for ranges of suitable markers (binning).

 ϕ_{s}

Our markers:

,
$$\widetilde{v}_{slip,y}$$





Some results













Filtered solid viscosity



Frdp300

Frdp150



11

0.6



Filtered solid pressure









Residual solid viscosity







Residual solid pressure







Main conclusion and future work



- The filtered parameters of concern were affected by particle Froude number at different extents, indicating that future constitutive models must account for particle Froude number.
- Before new constitutive models can be derived accounting for particle Froude number, further work is required to account for:
 # higher domain average gas Reynolds numbers
 # a variety of domain average solid volume fractions







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Thank you very much!