Adaptation of the vertical upflow phase map of Wirth to fluidized dense phase conveying of Geldart A powders and validation of the transition boundaries by Eulerian modelling with MFiX-TFM

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# Outline

# **Graphical Abstract**

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- Vertical upflow phase map
- Proposed (provisional) correlation for the Upper transport velocity (V<sub>tr2</sub>)
- Validation by Eulerian modelling
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  - Modelling options
  - Challenges in model validation
  - ➢ FF − DSU boundary at gross upflow
  - DSU dilute phase boundary
  - Packing limit
- Conclusions





# Fluidized Dense Phase Conveying (FDC)





# Fluidized dense phase conveying (FDC)

# Vertical up flow phase map





# Vertical upflow patterns of Geldart A powders

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#### Notes:

- For 'slightly cohesive' powders (HR < 1.25) which are easy to aerate and under well aerated conditions. They show clear pressure minima in the conveying characteristics.
- Impermeable plug flow may occur with slightly cohesive powders (HR < 1.25) in conventional FDC systems under inadequate aeration; sustained operation may require a suitable feeder.
- . Core-Annular Flow (CAF), Fast Fluidization (FF) and Dense Suspension Upflow (DSU) may occur at a same  $V_g$  but at different  $G_s$ : CAF at  $G_s < G_s^*$ , FF (LD-CFB) at  $G_s^* < G_s < G_s^d$  and DSU (HD-CFB) at  $G_s > G_s^d$



(Adapted from Kunii and Levenspiel, 1991 to include CAF and DSU)

Referenes for the flow patterns: Bi et al., 2000, 1993; Bi and Grace, 1999; Cocco et al., 2010; Kalman and Rawat, 2020; Klinzing et al., 2010; Kunii and Levenspiel, 1991; Liu et al., 1996; Loezos et al., 2002; Mills, 2004; Rabinovich and Kalman, 2011; Valverde, 2013; Wirth, 1988; Yerushalmi and Avidan, 1985





 $\Delta P/\Delta L$  – Pressure gradient;  $V_g$  – Superficial gas velocity;  $V_{mf}$  – Minimum fluidization velocity;  $V_{mb}$  – Minimum bubbling velocity;  $V_c$  – Onset of turbulent fluidization;  $V_{tr1}$  – Lower transport velocity,  $V_{tr2}$  – Upper transport

velocity;  $G_s$  – Solids flux;  $G_s^*$  – Saturation carrying capacity;  $G_s^d$  – Gross upflow flux;  $G_{s,tr2} - G_s$  corresponding to  $V_{tr2}$ ,

the threshold  $G_s$  for Dense Suspension Upflow;  $G_{s,tr1} - G_s$  corresponding to  $V_{tr1}$ 

### Vertical Up Flow Phase Map of Wirth adapted to FDC of Geldart A powders



#### • DSU and its transition boundaries are not broadly accepted [Grace et al. [1999]; Breault [2023]]

- Different oriteria have been proposed for the FF-DSU transition:
   1. Disappearance of the s-shaped axial c<sub>2</sub>profile. [Li and Kwauk (1980)]
   2. G<sub>2</sub> at which net upflow is attained in the dense annular region [Kim et al. (2004)]
- Recent experiment of Wang et al. (2022) [Fig.3 & 11] suggests that the FF DSU transition (defined in this project as the
  disappearance of the s-shaped axia e, profile) only occurs at gross upflow of solids in the entire riser [G<sup>2</sup><sub>7</sub>], with fittle
  backmixing. However, this requires validation, as the solids entering (although fluidized) at a lower velocity at the riser
  bottom could have formed a denser zone.
- The locus of pressure minima is broadly accepted as the boundary with dilute phase conveying at low  $G_s$  (hypically  $\varepsilon_s \cong 0.02$ ), however, its suitability for the high  $G_s$  of FDC requires validation.
- Resolved experimental data for the packing limit is scarce, and hence an hypothetical limit is proposed, considering that the dense locus (green) at  $V_g < V_{tr1}$  is unique and independent of  $G_s$ , and that the limit essentially lies where the available  $\Delta P / \Delta L$  balances the dominant losses under the dense conditions, static head of solids and drag.
- DSU and its transition boundaries are not broadly accepted [Grace et al.(1999); Breault(2023)].
- Different criteria have been proposed for the FF DSU transition :

   Disappearance of the s-shaped axial ε<sub>s</sub> profile. [Li and Kwauk (1980)]
   G<sub>s</sub> at which net upflow is attained in the dense annular region [Kim et al. (2004)]
- Recent experiment of Wang et al.(2022) (Fig.3 & 11) suggests that the FF DSU transition (defined in this project as the disappearance of the s-shaped axial  $\varepsilon_s$  profile) only occurs at gross upflow of solids in the entire riser ( $G_s^d$ ), with little backmixing. However, this requires validation, as the solids entering (although fluidized) at a lower velocity at the riser bottom could have formed a denser zone.
- The locus of pressure minima is broadly accepted as the boundary with dilute phase conveying at low  $G_s$  (typically  $\varepsilon_s \approx$  0.02), however, its suitability for the high  $G_s$  of FDC requires validation.
- Resolved experimental data for the packing limit is scarce, and hence an hypothetical limit is proposed, considering that the dense locus (green) at  $V_g < V_{tr1}$  is unique and independent of  $G_s$ , and that the limit essentially lies where the available  $\Delta P / \Delta L$  balances the dominant losses under the dense conditions, static head of solids and drag.

References for the phase map: Adapted from Wirth (1988), with inputs from Bi et al., 2000; Breault, 2023; Cocco et al., 2010; Kim et al., 2004; Li and Kwauk, 1980; Monazam and Shadle, 2011; Wang et al., 2022; Yerushalmi and Avidan, 1985

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Phase map of Wirth adapted for Geldart A powders

# Proposed correlation for the Upper transport velocity $(V_{tr2})$



# Proposed (provisional) correlation for the Upper transport velocity ( $V_{tr2}$ )



Note: A linear model ( $R^2 = 0.94$ ) or an exponential model ( $R^2 = 0.97$ ) may offer a better fit; however,  $Re_p - Ar$  correlations generally follow power law and further rigour is not attempted considering the uncertainty in the data.

### Correlation of Monazam & Shadle (2011):

Based on indirect measurement of  $V_{tr2}$  by column emptying times of solids (4 samples) that (tend to) group B.

$$Re_{p,tr2} = 3.118 \, Ar^{0.487}$$
 [1]

#### Proposed (provisional) correlation:

Based on 5 (sketchy) data points reported in the literature, for solids that clearly classify as group A, by direct measurement of  $V_{tr2}$  in CFB experiments at ambient conditions.

$$Re_{p,tr2} = 3.00 \, Ar^{0.80}$$
<sup>[2]</sup>

Provisional, considering the recognized uncertainties in the data, and as the dependence on riser diameter (**D**) and fraction of fines ( $P_{45}$ ) are not incorporated. However, it sufficiently demonstrates that  $V_{tr2}$  is significantly higher than predicted hitherto.

Powder	$d_p$	$\rho_s$	P <sub>45</sub>	D	L	L/D	Ar	V <sub>tr2</sub> (Reported), m/s		$Re_{P,tr2}$ $V_{tr2}$ (Predicted), n		cted), m/s	Reference
	μm	kg/m³		m	m			Reported	Used		Per [1]	Per [2]	
Catalyst	(30)	(1500)	Not reported	0.19	11.5	61	1.48	1.8	1.8	3.62	1.88	2.04	(Wirth, 1988) (Fig.7, pp.15)
FCC catalyst	55	729	Not reported	0.05	5.0	100	4.42	2.5 & 3.5	3.0	11.05	1.74	2.68	(Mori et al., 1992)
HFZ-20	49	1450	~ 0.32	0.15	8.5	56	6.22	> 4.1	5.0	16.41	2.31	3.95	(Yerushalmi and Avidan, 1985) (Fig.7.21 b, pp.259)
Alumina (fine)	54	3160	Not reported	0.09	8.0	89	17.75	> 5.0	6.0	21.70	3.50	8.29	(Li and Kwauk, 1980) (Fig.2, pp.540)
FCC catalyst	85	1500	~ 0.10	0.08	18.0	225	33.60	> 9.0	10.0	56.93	3.03	8.77	(Wang et al., 2022) (Fig.11 - II, pp.8)





# Proposed (provisional) correlation for the Upper transport velocity ( $V_{tr2}$ )

Recognized uncertainties in the  $V_{tr2}$  data points:

- Wirth (1988) had only reported  $V_t$  for the 'finely divided catalyst' powder;  $d_p$  has been estimated considering a  $\rho_s$  of 1500 kg/m<sup>3</sup>.
- Mori et al. (1992) had reported two different values for  $V_{tr2}$ , which is rather unique, affected by system design and operation; average of the values is used.
- Yerushalmi and Avidan (1985) had reported data for the HFZ-20 powder only up to a  $V_g$  of 4.1 m/s, at which the axial S-shaped  $\varepsilon_s$  profile is still very prominent; a  $V_{tr2}$  of 5 m/s is used.
- Li and Kwauk (1980) had reported data for the fine alumina only up to a  $V_g$  of ~ 5.0 m/s, and had projected a  $V_{tr2}$  of ~ 6.0 m/s.
- Wang et al. (2022) data for the FCC catalyst is limited to 9.0 m/s, at which they had approached (but not attained)  $V_{tr2}$ ; a  $V_{tr2}$  of 10 m/s is used.

### Differences in flow behaviour of Geldart A and B powders in CFB:

- Dependence of both  $V_{tr1}$  and  $V_{tr2}$  on  $G_s^*$  is intuitive.
- Experiments have shown that  $G_s^*$  increases with D for Geldart A but decreases for Geldart B; and in fact Breault et al. [Breault and Weber(2021); Breault et al.(2021)] have proposed separate correlations for groups A and B.
- While  $\varepsilon_s^*$  is ~ 0.01 for Geldart B, it increases with decreasing  $d_p$  and/or  $\rho_s$  to ~ 0.03 for Geldart A. [Bi et al. (1995)]
- $V_{tr1}$  Is a higher multiple of  $V_t$  for Geldart A than for Geldart B, also increasing with decreasing  $d_p$  and/or  $ho_s$ . [Bi et al. (1995)]

### Dependence of $V_{tr2}$ on riser diameter (D):

- Wang et al., (2022) approached (but not attained)  $V_{tr2}$  at a  $V_g$  of 9 m/s with a FCC powder ( $d_p$  (85 µm),  $\rho_s$  (1500 kg/m<sup>3</sup>) and  $P_{45}$  (~10 wt.%)) in a riser of D = 0.08 m.
- However, Issangya et al., (2023) observed very prominent axial S-shaped  $\varepsilon_s$  profile at a  $V_g$  of 12.2 m/s ( $G_s = 285 690 \text{ kg/m}^2\text{s}$ ) in a riser of D = 0.303 m, also with a FCC powder of similar  $d_p$  (79 µm),  $\rho_s$  (1490 kg/m<sup>3</sup>) and  $P_{45}$  (~8 wt.%).



### Proposed (provisional) correlation for the Upper transport velocity ( $V_{tr2}$ )



 $V_t$ ,  $V_{tr1}$  and  $V_{tr2}$  at ambient pressure relative to A-B boundary

 $V_t$  per Haider and Levenspiel (cited in Kunii & Levenspiel (1991))  $V_{tr1}$  per the correlation of Bi et al. (1995)  $V_{tr2}$  per the proposed (provisional) correlation.



Operating regions of FDC for finer and coarser group A powders

Published experimental data suggests that FDC for finer Geldart A powders, wherein the window of FF is narrow due to low  $V_{tr2}$  (and  $G_s^{tr2}$ ), operates at  $G_s > G_s^{tr1}$ . With increasing particle size  $(d_p)$  and / or density  $(\rho_s)$ , as the window of FF expands due to increasing  $V_{tr2}$ , the operation shifts to  $G_s^{tr2} < G_s < G_s^{tr1}$ , or even within the FF zone at  $G_s < G_s^{tr2}$ .



# Validation by Eulerian modelling with MFiX-TFM

- Powder characteristics
- Modelling options
- Challenges in model validation
- ➢ FF − DSU boundary at gross upflow
- DSU dilute phase boundary
- Packing limit



### Powder Characteristics

### Characteristics of the (hypothetical) powder at ambient conditions

Particle diameter	$d_p$	$\mu$ m	70	Monodisperse	
Particle density	$\rho_s$	Kg/m <sup>3</sup>	1400		
Archimedes number	Ar		17.5		
Void fraction loose packed	ε <sub>s,max</sub>		0.60	Maximum packing limit	
Interparticle restitution	c_e		0.95	Used for simulations	
Particle – wall restitution	e_w		0.89	Foerster et al.(1994), Drake(1991), (6mm cellulose acetate spheres)	
Angle of wall friction	Phi_w	0	11.86		
Angle of internal friction	Phi	0	30	McKeen & Pugsley (2003)	
Terminal settling velocity	V <sub>t</sub>	m / s	0.186		
Min. fluidization velocity	V <sub>mf</sub>	m / s	0.004	Kunii & Levenspiel (1991)	
Min. bubbling velocity	V <sub>mb</sub>	m / s	0.012		
Lower transport velocity	V <sub>tr1</sub>	m / s	1.4	Bi et al. (1995)	
Upper transport velocity	V <sub>tr2</sub>	m / s	6.3	Proposed (provisional) correlation	

Particles do not agglomerate or deform permanently; no electrostatic effects or liquid bridges.



#### Threshold for dominant van der Waals forces over gravity

Different criteria have been proposed for dominant van der Waals forces over gravity:

- (1)  $d_p = 55 \,\mu$ m, based on bed expansion [Loezos et al.(2002), Wang et al.(2011)] (2)  $Bo_q$  (granular bond number) = 1 [Valverde(2013)]

- (3) Ar = 16.5, based on pickup velocity [Kalman et al.(2005)]

Cohesion model is not considered essential for the monodisperse powder selected.

#### Validation by Eulerian modelling with MFiX-TFM

### Modelling options – Simulation conditions

Modelling options are largely in line with Balasubramanian et al.(2023)

Description	Nominal	Alternate		
MFiX version	22.3.1	22.3.1		
Viscous stress model	Simonin	Lun et al.(1984)		
Turbulence model	k-epsilon			
Frictional stress model	Princeton	Princeton		
Inlet & Outlet BC	MI & PO	MI & PO		
Wall BC for gas phase	NSW (Wall functions)	NSW		
Wall BC for solids phase	FSW / JJ-Mod	FSW		
Drag model	Di Felice	Di Felice		

#### Geometry and numerics Description Nominal grid Fine grid Riser (pipeline) diameter 0.04 0.04 D m Riser (pipeline) length 9.999 10.0 m L Grid 20 x 5000 40 x 9999 Cell size 0.001 m 0.002 Cell size in particle diameters 29 14 Maximum time step 1 x 10<sup>-5</sup> s 3 x 10<sup>-5</sup> Parallel processing DMP (1-16-1) DMP (1-8-1) Discretization Superbee Default Tolerances

### Cartesian 2D

#### (Reproduced from Balasubramanian et al. (2023))

Matches the wall profile (although in 2D).Qualitatively well predicts bed expansion, bubble rise, core- annular flow, clusters and streamers, etc.Lower computational cost for long pipe models (e.g., S- shaped profile).Offers rigorous wall BC.Disadvantages: Does not capture the inherent 3D nature of gas-solids flow. Numerical predictions can be affected by asymmetric flows, e.g., inlet and outlet configurations. (Li et al. 2014-I) May not simultaneously predict axial pressure profile and radial voidage accurately. (Li et al. 2014-II)		Advantages:							
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Optimum underrelaxation factors that prevent backflow with pressure outlet BC (Reproduced from Balasubramanian et al. (2023))

For 2D Cartesian grid simulations #								
Under r	elaxation factor	Default	Selected	Remarks				
UR_FAC(1)	Gas pressure	0.8	0.8	Improves stability				
UR_FAC(2)	EP_s	0.5	1.0					
UR_FAC(3&5)	U&W-Momentums	0.5	0.5	Improves stability				
UR_FAC(4)	V-Momentums	0.5	<mark>1.0</mark>	Set at 1.0 to avoid backflow at no / low solids loadings				
UR_FAC(8)	Granular temperature	0.5	0.5	~ 1 in 3 runs diverges without underrelaxation				
UR_FAC(9)	k-e	0.8	<mark>1.0</mark>	Set at 1.0 to avoid backflow at no / low solids loadings				
UR_F_GS	Drag	1.0	0.0	Improves stability at ~ 3% increase in wall time				

# - May not be suitable for other coordinate systems or simulation conditions

Validation by Eulerian modelling with MFiX-TFM

Reproduced from Balasubramanian et al. (2023)

### Modelling options – Model scheme



particle relaxation time scales; Simonin's turbulence model and the GKT are recovered at the dilute and dense limits, respectively.

 $au_s$  - Dissipation time scale, s

 $au_{gs}^{x}$  - Particle relaxation time, s

 $au_s^c$  - Collisional time scale, s

Reference: Balzer et al. (1996) Benyahia et al. (2005) & (2007), Srivastava & Sundaresan (2003)











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Validation by Eulerian modelling with MFiX-TFM

### Validation: DSU – Dilute phase boundary and packing limit



Solids flux = 2000 kg / m<sup>2</sup> s

Conditions at the feed zone:  $V_g = 3 \text{ m/s}$  at 4 bar



• For the conditions simulated, pressure minimum is not attained at the ambient receiver.

• Trends in granular temperature and slip velocity are probable aternate variables to perceive the transition.

• As precited by Granular Kinetic Theory, gravity dominates the pressure drop at the packing limit and drag is insignificant (due to the very low slip velocity).

• Contribution of other pressure losses (drag, etc.) increase from ~ 2% at the packing limit to > 20% at the ambient receiver.

## Conclusions

 Vertical upflow phase map of Wirth adapted to locate Fluidized Dense Phase Conveying of Geldart A powders. The map demarcates the boundaries of Dense Suspension Upflow with Fast Fluidization, dilute phase conveying and the packing limit.



• Proposed a provisional correlation for the Upper transport velocity ( $V_{tr2}$ ) of group A powders, based on a limited (sketchy) data set, without accounting for its dependence on riser diameter (D) and fines fraction ( $P_{45}$ ). The correlation sufficiently demonstrates that is significantly higher than hitherto predicted.

$$Re_{p,tr2} = 3.00 Ar^{0.80}$$

- Highlighted the challenges faced in validating the MFiX-TFM model at high solids flux: low solids concentration near the wall and lower slip velocities than reported in CFB experiments.
- Demonstrated (based on the experimental results of Wang et al.(2022) and Eulerian modelling) that the transition from Fast Fluidization to Dense Suspension Upflow, defined as the disappearance of the S-shaped axial profile only occurs at gross upflow of solids in the entire riser.
- For the conditions simulated, transition to dilute phase conveying is not perceived at the ambient receiver based on the pressure gradient.
- Packing limit, as predicted by the Granular Kinetic Theory is largely due to static head of solids (as the predicted slip velocity is very low).

Areas identified for further research:

- 1. CFB experiments to measure  $V_{tr2}$  of group A powders, incorporating various riser diameters and fines fractions, for a robust correlation.
- 2. Validation of the TFM model for high solids flux applications.





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